การแปลงซีกัล-บาร์กมันน์บนทรงกลม



## SEGAL-BARGMANN TRANSFORM ON SPHERES



ISBN 974-53-1953-8

Thesis Title
By Mr. Boonyong Sriponpaew
Field of Study
Thesis Advisor
Mathematics

SEGAL-BARGMANN TRANSFORM ON SPHERES

Assistant Professor Wicharn Lewkeeratiyutkul, Ph.D.

Accepted by the Faculty of Science, Chulalongkorn University in Partial Fulfillment of the Requirements for the Master's Degree

Deputy Dean for Administrative Affairs, Acting Dean, The Faculty of Science (Associate Professor Tharapong Vitidsant, Ph.D.)

Thesis Committee

(Assistant Professor Imchit Termwuttipong, Ph.D.)

(Assistant Professor Wicharn Lewkeeratiyutkul, Ph.D.)

................................ Member
(Songkiat Sumetkijakan, Ph.D.)

บุญยงค์ ศรีพลแผ้ว : การแปลงซีกัล-บาร์กมันน์บนทรงกลม(SEGAL-BARGMANN ON SPHERES) อ. ที่ปรึกษา : ผศ. ดรวิชาญ ลิ่วกีรติยุตกุล, 30 หน้า. ISBN 974-53-1953-8

เราศึกษาการแปลงซีกัล-บาร์กมันน์บนทรงกลม $d$ มิติ $S^{d}$ ซึ่งส่งค่าจาก $L^{2}\left(S^{d}\right)$ ไปยัง ฟังก์ชันโฮ โลโมฟิกบน ทรงกลมเชิงซ้อน $\mathrm{S}_{\mathrm{C}}^{\mathrm{d}}$ ในการวิจัยครั้งนี้ เราได้พิสูจน์ว่าการแปลงซีกัล-บาร์กมันน์เป็นฟังก์ชัน สมมติจาก $\mathrm{L}^{2}\left(\mathrm{~S}^{\mathrm{d}}\right)$ ไปทั่วถึง ปริภูมิของฟังก์ชันโฮโลโมฟิกบน $\mathrm{S}_{\mathrm{C}}^{\mathrm{d}}$ ซึ่งสามารถอินทิเกรตกำลังสองเทียบกับ เมเชอร์บางชนิดได้


## สถาบันวิทยบริการ

## จุฬาลงกรณ์มหาวิทยาลัย

ภาควิชา คณิตศาสตร์
สาขาวิชา คณิตศาสตร์
ปีการศึกษา 2547

ลายมือชื่อนิสิต.
ลายมือชื่ออาจารย์ที่ทรึกษา.
$\qquad$
\# \# 4572353023 : MAJOR MATHEMATICS
KEY WORDS : SEGAL-BARGMANN / SPHERE

## BOONYONG SRIPONPAEW : SEGAL-BARGMANN TRANSFORM ON

SPHERES : THESIS ADVISOR : ASST. PROF. WICHARN
LEWKEERATIYUTKUL, Ph.D. 30 pp. ISBN 974-53-1953-8

We study the Segal-Bargmann transform on the $d$-dimensional sphere $S^{d}$, mapping $L^{2}\left(S^{d}\right)$ into holomorphic functions on the complexification $S_{\mathbb{C}}^{d}$. In this work, we prove that the Segal-Bargmann transform is an isometry from $L^{2}\left(S^{d}\right)$ onto the space of holomorphic functions on $S_{\mathbb{C}}^{d}$ which are square integrable with respect to a certain measure.



Department Mathematics
Field of study Mathematics

Student's signature
Advisor's signature
$\qquad$
$\qquad$

Academic year 2004

## Acknowledgments

First of all, I would like to express great appreciation to my advisor Assist. Prof. Dr. Wicharn Lewkeeratiyutkul for his supreme knowledge and guidance. All the time, he has always taught me not only academic study but also the way of my life. My special thanks are given to the thesis committee: Assistant Professor Imchit Termwuttipong, Dr. Songkiat Sumetkijakan, Dr. Paolo Bertozzini. They kindly read the thesis and offered many useful suggestions on the suitability of the thesis content. I am also grateful to the Development and Promotion of Science and Technology Talents Project (DPST) for the funding that I received for studying and doing research. For 7 years of my study at the Department of Mathematics, Chulalongkorn University, I feel very thankful to my department and all of my teachers who have taught me knowledge and skills. Thanks also goes to all of my friends who have always encouraged me until I finished the course of study. Finally, I feel very grateful to my parents, who have brought me up, stood by me and given me extremely valuableadvice. 6100190 a

## Table of Contents

Abstract in Thai ..... iv
Abstract in English ..... V
Acknowledgments ..... vi
Table of Contents ..... vii
1 Introduction ..... 1
2 Complex sphere ..... 3
3 Spherical harmonics ..... 7
4 Laplacian on a sphere ..... 13
5 Segal-Bargmann transform ..... 17
Bibliography ..... 29
Vita ..... 30
สถาบันวิทยบริการ จุฬาลงกรณ์มหาวิทยาลัย

## Chapter 1

## Introduction

The Segal-Bargmann transform is a transform which is widely studied by physicists. It is used for describing the wave-particle phenomena in quantum field theory. The Segal-Bargmann transform for $\mathbb{R}^{d}$ is a unitary map $C_{t}: L^{2}\left(\mathbb{R}^{d}\right) \rightarrow \mathcal{H} L^{2}\left(\mathbb{C}^{d}, \nu_{t}\right)$ defined by

$$
C_{t} f(z)=\int_{\mathbb{R}^{d}}(2 \pi t)^{-\frac{d}{2}} e^{-(z-x)^{2} / 2 t} f(x) d x, \quad z \in \mathbb{C}^{d}
$$

where $(z-x)^{2}=\left(z_{1}-x_{1}\right)^{2}+\left(z_{2}-x_{2}\right)^{2}+\cdots+\left(z_{d}-x_{d}\right)^{2}$ and $\mathcal{H} L^{2}\left(\mathbb{C}^{d}, \nu_{t}\right)$ denotes the space of holomorphic functions that are square integrable with respect to the measure

$$
66 \nu_{t}(z)=(2 \pi t)^{-\frac{d}{2}} e^{-|\operatorname{Im} z|^{2} / 2 t} \cdot \square \approx
$$

There are a lot of generalizations of the Segal-Bargmann transform on $\mathbb{R}^{d}$ to more general settings. In 1993, Hall $([3])$ has obtained a generalization of the SegalBargmann transform on a compact Lie group which is geometric in nature and keeps more of a structure of the original Segal-Bargmann transform. In his work, the space $\mathbb{R}^{d}$ is replaced by a connected compact Lie group $K$ and $\mathbb{C}^{d}$ is replaced
by the complexification $K_{\mathbb{C}}$ of $K$. Later, the Segal-Bargmann transform was also extended by Stenzel ([7]) to the case of compact symmetric spaces. His proof relies on heavy machinery in theory of symmetric spaces. In this thesis we study the SegalBargmann transform on the $d$-dimensional sphere $S^{d}$, which is a special case of a compact symmetric space. However, we define and prove everything explicitly using only elementary methods. Indeed, we follow the outline already given by Hall and Mitchell ([4]). We provide necessary backgrounds and give the proofs in complete detail.


## Chapter 2

## Complex sphere

The $d$-dimensional sphere is the subset of $\mathbb{R}^{d+1}$ given by

$$
S^{d}=\left\{x \in \mathbb{R}^{d+1}+x_{1}^{2}+\cdots+x_{d+1}^{2}=1\right\}
$$

We naturally define the complexified sphere $S_{\mathbb{C}}^{d}$ to be

$$
S_{\mathbb{C}}^{d}=\left\{z \in \mathbb{C}^{d+1} \mid z_{1}^{2}+\cdots+z_{d+1}^{2}=1\right\}
$$

Then $S_{\mathbb{C}}^{d}$ is a $d$-dimensional complex manifold. We will show that we can identify $S_{\mathbb{C}}^{d}$ with the cotangent bundle $T^{*}\left(S^{d}\right)$ of $S^{d}$, which is defined by

$$
T^{*}\left(S^{d}\right)=\left\{(\mathbf{x}, \mathbf{p}) \in \mathbb{R}^{d+1} \times \mathbb{R}^{d+1}|\| \mathbf{x}|=1, \mathbf{x} \cdot \mathbf{p}=0\right\}
$$

Notice that $T^{*}\left(S^{d}\right)$ is a $2 d$-dimensional real manifold. If we view $S_{\mathbb{C}}^{d}$ as a $2 d$ dimensional real manifold, then we can identify these two manifolds together via the following map

$$
\mathbf{a}(\mathbf{x}, \mathbf{p})=(\cosh p) \mathbf{x}+i \frac{\sinh p}{p} \mathbf{p}
$$

where $(\mathbf{x}, \mathbf{p}) \in T^{*}\left(S^{d}\right)$ and $p=|\mathbf{p}|$. Since $\lim _{p \rightarrow 0} \frac{\sinh p}{p}=1$, it is well-defined when $p=0$. First note that

$$
\begin{aligned}
\mathbf{a}(\mathbf{x}, \mathbf{p}) \cdot \mathbf{a}(\mathbf{x}, \mathbf{p}) & =\sum_{k=1}^{d+1}\left(x_{k}^{2} \cosh ^{2} p-\frac{p_{k}^{2}}{p^{2}} \sinh ^{2} p\right)+i\left(\frac{2 p_{k} x_{k}}{p} \sinh p \cosh p\right) \\
& =\cosh ^{2} p-\sinh ^{2} p \\
& =1
\end{aligned}
$$

This shows that a maps $T^{*}\left(S^{d}\right)$ into $S_{\mathbb{C}}^{d}$. To prove that it is injective, let $(\mathbf{x}, \mathbf{p})$, $(\mathbf{y}, \mathbf{q}) \in T^{*}\left(S^{d}\right)$ be such that

$$
\mathbf{a}(\mathbf{x}, \mathbf{p})=\mathbf{a}(\mathbf{y}, \mathbf{q})
$$

Then

$$
x_{k} \cosh p+i p_{k} \frac{\sinh p}{p}=y_{k} \cosh q+i q_{k} \frac{\sinh q}{q}
$$

i.e.,

$$
x_{k} \cosh p=y_{k} \cosh q \quad \text { and } \quad p_{k} \frac{\sinh p}{p}=q_{k} \frac{\sinh q}{q}
$$

for $k=1, \ldots, d+1$. This implies that

$$
\mathbf{x}^{2} \cosh ^{2} p=\mathbf{y}^{2} \cosh ^{2} q \quad \text { and } \quad \sinh ^{2} p=\sinh ^{2} q
$$

Consequently $\sinh p= \pm \sinh q$. Since $p, q \geq 0$, it follows that $\sinh p=\sinh q$ and that $p=q$. Hence $x_{k}=y_{k}$ and $p_{k}=q_{k}$ for every $k$.
Next, we show that $\mathbf{a}$ is surjective. Let $\mathbf{z} \in S_{\mathbb{C}}^{d}$ and write $\mathbf{z}=\mathbf{r}+\mathbf{i}$, where $\mathbf{r}, \mathbf{s} \in \mathbb{R}^{d+1}$. Choose

$$
\begin{equation*}
\mathbf{p}=\frac{\sinh ^{-1}|\mathbf{s}|}{|\mathbf{s}|} \mathbf{s} \quad \text { and } \quad \mathbf{x}=\frac{\mathbf{r}}{\cosh p} . \tag{2.1}
\end{equation*}
$$

Then

$$
\frac{\sinh p}{p} \mathbf{p}=\frac{\sinh \left(\sinh ^{-1}|\mathbf{s}|\right)}{\sinh ^{-1}|\mathbf{s}|} \sinh ^{-1}|\mathbf{s}| \frac{\mathbf{s}}{|\mathbf{s}|}=\mathbf{s} \quad \text { and } \quad \cosh p \mathbf{x}=\mathbf{r} .
$$

Hence $\mathbf{a}(\mathbf{x}, \mathbf{p})=\mathbf{z}$. From (2.1), we can see that

$$
\mathbf{a}^{-1}(\mathbf{r}+i \mathbf{s})=\left(\frac{\mathbf{r}}{\cosh p}, \frac{\sinh ^{-1}|\mathbf{s}|}{|\mathbf{s}|} \mathbf{s}\right) .
$$

It is clear that $\mathbf{a}$ and $\mathbf{a}^{-1}$ are smooth functions. Hence $\mathbf{a}$ is a diffeomorphism.
Next we will show that the smooth manifold $S^{d}$ is diffeomorphic to the homogeneous manifold $S O(d+1) / S O(d)$. Let us recall some definitions and theorems about homogeneous manifolds first.

Definition 2.1. Let $\eta: G \times M \rightarrow M$ be an action of $G$ on $M$ on the left. As usual, we write

$$
\eta(g, m)=g \cdot m .
$$

The action is called transitive if whenever $m$ and $n$ belong to $M$ there exists a $g$ in $G$ such that $g \cdot m=n$. For $x_{0} \in M$, the set

$$
G_{x_{0}}=\left\{g \in G \mid g \cdot x_{0}=x_{0}\right\}
$$

forms a closed subgroup of $G$ called the isotropy group at $x_{0}$.
If $G$ is a Lie group and $H$ is a closed subgroup of $G$, then we can define a differentiable structure on the quotient space $G / H$ so that it is a smooth manifold, called a homogeneous manifold. Moreover, there is a natural transtive left-action of $G$ on $G / H$. Conversely, if $M$ is a smooth manifold and there is atransitive leftaction by a Lie group $G$ on $M$, then $M$ can be identified with the quotient manifold $G / G_{x_{0}}$, where $x_{0}$ is a point in $M$. This is summarized in the following theorem.

Theorem 2.2 ([9], Theorem 3.62). Let $\eta: G \times M \rightarrow M$ be a transitive leftaction of the Lie group $G$ on the manifold $M$. Let $x_{0} \in M$, and let $H$ be the isotropy group at $x_{0}$. Define a mapping $\beta: G / H \rightarrow M$ by $\beta(g H)=g \cdot x_{0}$. Then $\beta$ is a diffeomorphism.

Let $\mathbb{F}$ be the field $\mathbb{R}$ and $\mathbb{C}$. For $d \geq 1$, we define the special orthogonal group $S O(d, \mathbb{F})$ to be the set of $d \times d$ matrices $A$ such that $A \cdot A^{t}=I$ and $\operatorname{det} A=1$. Equivalently, $S O(d, \mathbb{F})$ is the set of $d \times d$ matrices $A$ such that $\operatorname{det} A=1$ and $[A x, A y]=[x, y]$ for all $x, y \in \mathbb{F}$, where $[x, y]=\sum_{i=1}^{d} x_{i} y_{i}$ for $x, y \in \mathbb{F}^{d}$.

In case $\mathbb{F}=\mathbb{R}$, we may simply write $S O(d, \mathbb{R})$ as $S O(d)$.
Proposition 2.3. The manifold $S^{d}$ is diffeomorphic to $S O(d+1) / S O(d)$.

Proof. Let $\left\{e_{i} \mid i=1, \ldots, d+1\right\}$ be the canonical basis of $\mathbb{R}^{d+1}$ where $e_{i}$ is the $d+1$-tuple consisting of all zeroes except for a 1 in the $i$-th spot. Define an action $\eta: S O(d+1) \times S^{d} \rightarrow S^{d}$ by multiplying $A \in S O(d+1)$ to a vector in $S^{d}:$

$$
\eta(A, x)=A x \text {. }
$$

It is obvious that this action is transitive and the isotropy group at $e_{d+1}$ is the set of matrices in $S O(d+1)$ of the form

The matrix $B$ occuring in this subgroup is precisely the matrix in $S O(d)$. Hence, we identify this isotropy group with $S O(d)$. It follows from Theorem 2.2 that the homogeneous manifold $S O(d+1) / S O(d)$ is diffeomorphic to $S^{d}$.

Similarly, we can show that $S_{\mathbb{C}}^{d}$ is diffeomorphic to $S O(d+1, \mathbb{C}) / S O(d, \mathbb{C})$ as complex manifolds.


## Chapter 3

## Spherical harmonics

We would like to represent a function defined on the surface of the unit sphere by an expansion similar to a Fourier series by a class of functions called the spherical harmonics. For more details, a reader is referred to [1] and [8].

A function $f$ defined on $\mathbb{R}^{d}$ is said to be homogeneous of degree $k$ if

$$
f(a x)=a^{k} f(x) \quad \text { for any } x \in \mathbb{R}^{d} \text { and any } a>0
$$

Let $\mathcal{P}_{k}\left(\mathbb{R}^{d}\right)$ be the set of all homogeneous polynomials of degree $k$ on $\mathbb{R}^{d}$. If $P \in$ $\mathcal{P}_{k}\left(\mathbb{R}^{d}\right)$, then it can be written in the form

$$
66 \rightarrow 9(x)=\sum_{|\alpha|=k} c_{\alpha} x^{\alpha}
$$

where $\alpha$ denotes a $d$-tuple $\left(\alpha_{1}, \alpha_{2}, \ldots, \alpha_{d}\right)$ of nonnegative integers, $|\alpha|=\alpha_{1}+\alpha_{2}+$
$\cdots+\alpha_{d}, x^{\alpha}=x_{1}^{\alpha_{1}} \ldots x_{d}^{\alpha_{d}}$ and $c_{\alpha} \in \mathbb{C}$.
It is clear that the set of monomials $\left\{x^{\alpha}:|\alpha|=k\right\}$ is a basis for this space. Then $\operatorname{dim} \mathcal{P}_{k}\left(\mathbb{R}^{d}\right)$ equals the number of distinct multi-indices $\alpha=\left(\alpha_{1}, \alpha_{2}, \ldots, \alpha_{d}\right)$ with
$|\alpha|=k$. Hence

$$
\operatorname{dim} \mathcal{P}_{k}\left(\mathbb{R}^{d}\right)=\binom{d+k-1}{d-1}
$$

We introduce an inner product $\langle P, Q\rangle$ on $\mathcal{P}_{k}$ by letting $\langle P, Q\rangle=P(D) \bar{Q}$ for all $P$, $Q$ in $\mathcal{P}_{k}$, where $P(D)$ is the differential operator in which we replace $x_{1}^{\alpha_{1}} x_{2}^{\alpha_{2}} \ldots x_{d}^{\alpha_{d}}$ by $\frac{\partial^{\alpha_{1}+\cdots+\alpha_{d}}}{\partial x_{1}^{\alpha_{1}} \partial x_{2}^{\alpha_{2}} \cdots \partial x_{d}^{\alpha_{d}}}$. Since $P$ and $Q$ are homogeneous polynomials of the same degree, $\langle P, Q\rangle$ is scalar-valued. It is clearly linear in the first variable, conjugate linear in the second variable and hermitian symmetric. To verify that it is an inner product, it is enough to show that $\langle P, P\rangle \geq 0$, with equality only if $P=0$. If $\alpha \neq \beta$, then

$$
\left(\frac{\partial^{\alpha_{1}}}{\partial x_{1}^{\alpha_{1}}} \frac{\partial^{\alpha_{2}}}{\partial x_{2}^{\alpha_{2}}} \ldots \frac{\partial^{\alpha_{d}}}{\partial x_{d}^{\alpha_{d}}}\right) x_{1}^{\beta_{1}} x_{2}^{\beta_{2}} \ldots x_{d}^{\beta_{d}}=0 .
$$

When $\alpha=\beta$, this derivative equals $\alpha_{1}!\alpha_{2}!\ldots \alpha_{d}!=\alpha!$. Consequently, if $P(x)=$ $\sum_{|\alpha|=k} c_{\alpha} x^{\alpha}$, then $\langle P, P\rangle=\sum_{|\alpha|=k}\left|c_{\alpha}\right|^{2} \alpha!$. But this last expression vanishes if and only if all the coefficients $c_{\alpha}$ are 0 .

Theorem 3.1. If $P \in \mathcal{P}_{k}$, then

$$
P(x)=P_{0}(x)+|x|^{2} P_{1}(x)+\cdots+|x|^{2 l} P_{l}(x),
$$

where $P_{j}$ is a homogeneous harmonic polynomial of degree $k-2 j$, for $j=0,1, \ldots, l$. Proof. Any polynomial of degree less than 2 isharmonic. Thus we may assume that $k \geq 2$. Consider the linear mapping $\varphi: \mathcal{P}_{k} \rightarrow \mathcal{P}_{k-2}$ defined by letting $\varphi(P)=\Delta P$ for $P_{0} \in \mathcal{P}_{k}$, where $\Delta$ is the Laplace operator on $\mathbb{R}^{d}$. We first show that $\varphi$ maps $\mathcal{P}_{k}$ ont8 $\mathcal{P}_{k-2}$. If this were not the case we could find a nonzero $Q \in \mathcal{P}_{k-2}$ that is orthogonal to Range $(\varphi)$. That is

$$
\overline{\langle\Delta P, Q\rangle}=\langle Q, \Delta P\rangle=0, \text { for all } P \in \mathcal{P}_{k} .
$$

In particular, this must be true for $P(x)=|x|^{2} Q(x)$. Thus

$$
0=\langle Q, \Delta P\rangle=Q(D) \overline{\Delta P}=\Delta Q(D) \bar{P}=P(D) \bar{P}=\langle P, P\rangle .
$$

But this is impossible since $P \neq 0$.
Let $\mathcal{A}_{j} \subseteq \mathcal{P}_{j}, j \geq 2$, be the class of all harmonic polynomials in $\mathcal{P}_{j}$ and $\mathcal{B}_{j}=$ $|x|^{2} \mathcal{P}_{j-2}$.

We claim that $\mathcal{P}_{j}$ is the orthogonal direct sum of $\mathcal{A}_{j}$ and $\mathcal{B}_{j}$.

$$
\begin{aligned}
\left.\left.\langle | x\right|^{2} Q, P\right\rangle \text { for all } Q \in \mathcal{P}_{j-2} & \Leftrightarrow Q(D) \Delta \bar{P}=0 \text { for all } Q \in \mathcal{P}_{j-2} \\
& \Leftrightarrow\langle Q, \Delta P\rangle=0 \text { for all } Q \in \mathcal{P}_{j-2} \\
& \Leftrightarrow \Delta P=0 .
\end{aligned}
$$

In particular, for $j=k$ and $P \in \mathcal{P}_{k}$ we have $P(x)=P_{0}(x)+|x|^{2} Q(x)$ with $P_{0}$ harmonic and $Q \in \mathcal{P}_{k-2}$. It is clear that the desired statement follows by induction.

Corollary 3.2. The restriction to the unit sphere $S^{d-1}$ of any polynomial of $d$ variables is a sum of restrictions to $S^{d-1}$ of harmonic polynomials.

The restriction to the unit sphere $S^{d-1}$ of a homogeneous harmonic polynomial of degree $k$ is called a spherical harmonic of degree $k$. We let $\mathcal{H}_{k}$ denote the space of spherical harmonics of degree $k$.

Let $\varphi: \mathcal{A}_{k} \rightarrow \mathcal{H}_{k}$ be defined by 9 ¢

## จฬาลงกรใซึตตาวิทยาลัย

It is evident that this map has a trivial kernel. If $Y \in \mathcal{H}_{k}$, we can choose $P(x)=$ $x^{k} Y(x /|x|)$ for $x \neq 0$. Then $\varphi$ is an isomorphism of $\mathcal{A}_{k}$ onto $\mathcal{H}_{k}$. Hence,

$$
\begin{equation*}
\operatorname{dim} \mathcal{H}_{k}=\operatorname{dim} \mathcal{A}_{k}=\operatorname{dim} \mathcal{P}_{k}-\operatorname{dim} \mathcal{P}_{k-2}=\frac{(d+2 k-2)}{k}\binom{d+k-3}{k-1} . \tag{3.1}
\end{equation*}
$$

To prove the next proposition, let us recall Green's theorem.

Theorem 3.3 (Green's theorem). Let $u, v \in C^{2}(\bar{U})$, where $U$ is bounded subset of $\mathbb{R}^{d}$. Then

$$
\int_{U}(u \Delta v-v \Delta u) d V=\int_{\partial U}\left(u \partial_{n} v-v \partial_{n} u\right) d s
$$

where $\partial_{n}$ denotes differentiation with respect to the outward unit normal vector.
Proposition 3.4. If $Y_{k}$ and $Y_{l}$ are spherical harmonics of degree $k$ and $l$, with $k \neq l$, then

$$
Y_{k}(x) Y_{l}(x) d x=0
$$

Proof. By Green's theorem,

$$
\int_{S^{d-1}} Y_{k} \partial_{n} Y_{l}-Y_{l} \partial_{n} Y_{k} d s=\int_{|x| \leq 1}\left(Y_{k} \Delta Y_{l}-Y_{l} \Delta Y_{k}\right) d x=0 .
$$

But then for each $x \in S^{d-1}$,


Similarly, $\partial_{n} Y_{l}=l Y_{l}$. Thus

Since $a \neq k$, the last integral vanishes, as desired.


Define $L^{2}\left(S^{d-1}\right)$ to be the Hilbert space of square-integrable functions on the (d-1)-dimensional sphere $S^{d-1}$ with respect to surface measure $d x$. Then each $\mathcal{H}_{k}$ is a subspace of $L^{2}\left(S^{d-1}\right)$. Moreover, we have

Theorem 3.5. $L^{2}\left(S^{d-1}\right)=\bigoplus_{k=0}^{\infty} \mathcal{H}_{k}$.
Proof. $L^{2}\left(S^{d-1}\right)=\bigoplus_{k=0}^{\infty} \mathcal{H}_{k}$ is true when the following three conditions are satisfied:

1. $\mathcal{H}_{k}$ is a closed subspace of $L^{2}\left(S^{d-1}\right)$ for each $k$;
2. $\mathcal{H}_{k}$ is orthogonal to $\mathcal{H}_{l}$ for $k \neq l$;
3. For every $f \in L^{2}\left(S^{d-1}\right)$, there exists a sequence $\left(f_{m}\right)$, where $f_{m} \in \mathcal{H}_{m}$ for each $m$, such that

$$
f=f_{0}+f_{1}+\ldots,
$$

where the sum converges in the norm of $L^{2}\left(S^{d-1}\right)$.
Condition 1 above holds because each $\mathcal{H}_{k}$ is finite dimensional and hence is closed in $L^{2}\left(S^{d-1}\right)$. Condition 2 follows from the Proposition 3.4.

To prove condition 3, we only need to show that the linear span of $\cup_{k=0}^{\infty} \mathcal{H}_{k}$ is dense in $L^{2}\left(S^{d-1}\right)$. As we have already noted from the Corollary 3.2 that if $P$ is a polynomial on $\mathbb{R}^{d}$, then $\left.P\right|_{S^{d-1}}$ can be written as a finite sum of elements of $\cup_{k=0}^{\infty} \mathcal{H}_{k}$. By the Stone-Weierstrass theorem, the set of restrictions $\left.P\right|_{S^{d-1}}$, as $P$ ranges over all polynomials on $\mathbb{R}^{d}$, is dense in $C\left(S^{d-1}\right)$ with respect to the supremum norm. Because $C\left(S^{d-1}\right)$ is dense in $L^{2}\left(S^{d-1}\right)$ and the $L^{2}$-norm is less than or equal to the supremum norm on $S^{d-1}$, this implies that the linear span of $\cup_{k=0}^{\infty} \mathcal{H}_{k}$ is dense in


If $\left\{Y_{\mathrm{P}, k}, \cdots, Y_{N_{k}, k}\right\}$ is an orthonormal basis of $\mathcal{H}_{k}$, then it follows from Theorem 3.5 that the collection $\cup_{k=0}^{\infty}\left\{Y_{1, k}, \ldots, Y_{N_{k}, k}\right\}$ is an orthonormal basis of $L^{2}\left(S^{d-1}\right)$. Thus, if $f \in L^{2}\left(S^{d-1}\right)$ then there exists a unique representation,

$$
f=\sum_{k=0}^{\infty} \sum_{n=1}^{N_{k}}\left\langle f, Y_{n, k}\right\rangle Y_{n, k}
$$

where the series on the right converges to $f$ in the $L^{2}$ norm. Let us fix a point $x$ in $S_{\mathbb{C}}^{d-1}$ and consider the linear functional $L$ on $\mathcal{H}_{k}$ that assigns to each $Y$ in $\mathcal{H}_{k}$ the value $Y(x)$. By Riesz representation theorem, there exists a unique spherical harmonic $Z_{x}^{(k)}$ such that

$$
L(Y)=Y(x)=\int_{S^{d}-1} Y(t) Z_{x}^{(k)}(t) d t
$$

for any $Y \in \mathcal{H}_{k}$. This function $Z_{x}^{(k)}$ is called the zonal harmonic of degree $k$ with pole $x$.

Lemma 3.6. Let $\left\{Y_{1}, Y_{2}, \ldots, Y_{N_{k}}\right\}$ be an orthonormal basis of $\mathcal{H}_{k}$. Then
(i) $Z_{x}^{(k)}(t)=\sum_{m=1}^{N_{k}} \overline{Y_{m}(x)} Y_{m}(t)$;
(ii) If $\rho$ is a rotation, then $Z_{\rho x}^{(k)}(\rho t)=Z_{x}^{(k)}(t)$;
(iii) $\sum_{m=1}^{N_{k}}\left|Y_{m}(x)\right|^{2}=\frac{N_{k}}{C_{d-1}}$ where $C_{d-1}=\int_{S^{d-1}} d x$ is the total volume of $S^{d-1}$;
(iv) If $Y \in \mathcal{H}_{k}$, then $|Y(x)|^{2} \leq \frac{N_{k}}{C_{d-1}}\|Y\|_{2}^{2}$, where $N_{k}=\operatorname{dim} \mathcal{H}_{k}$.

Proof. Since $\left\{Y_{1}, Y_{2}, \ldots, Y_{N_{k}}\right\}$ is an orthonormal basis of $\mathcal{H}_{k}$,

$$
Z_{x}^{(k)}=\sum_{m=1}^{N_{k}}\left\langle Z_{x}^{(k)}, Y_{m}\right\rangle Y_{m}
$$

But by the defining property of zonal harmonics, $\stackrel{\square}{6}$

Then $Z_{x}^{(k)}(t)=\sum_{m=1}^{N_{k}} \overline{Y_{m}(x)} Y_{m}(t)$. To verify (ii), let $w=\rho t$. We have

$$
\int_{S^{d-1}} Z_{\rho x}^{(k)}(\rho t) Y(t) d t=\int_{S^{d-1}} Z_{\rho x}^{(k)}(w) Y\left(\rho^{-1} w\right) d w=Y\left(\rho^{-1} \rho x\right)=Y(x)
$$

By the uniqueness of zonal harmonic, we have $Z_{\rho x}^{(k)}(\rho t)=Z_{x}^{(k)}(t)$. To verify (iii), suppose $x_{1}$ and $x_{2}$ are in $S^{d-1}$. We can find a rotation $\rho$ in $S O(d)$ such that $\rho x_{1}=x_{2}$. Then

$$
Z_{x_{2}}^{(k)}\left(x_{2}\right)=Z_{x_{1}}^{(k)}\left(x_{1}\right) .
$$

Consequently, $Z_{x}^{(k)}(x)$ is a constant, say c. From (i), we have $c=Z_{x}^{(k)}(x)=$ $\sum_{m=1}^{N_{k}}\left|Y_{m}(x)\right|^{2}$. Since $\left\|Y_{m}\right\|_{2}=1$ for all $m$,

$$
N_{k}=\sum_{m=1}^{N_{k}} \int_{S^{d-1}}\left|Y_{m}(x)\right|^{2} d x=\int_{S^{d-1}} \sum_{m=1}^{N_{k}}\left|Y_{m}(x)\right|^{2} d x=\int_{S^{d-1}} c d x=c C_{d-1}
$$

Hence $\sum_{m=1}^{N_{k}}\left|Y_{m}(x)\right|^{2}=c=N_{k} / C_{d-1}$, so (iii) follows. From (i) it implies that

$$
\left\|Z_{x}^{(k)}\right\|_{2}^{2}=\int_{S^{d-1}}\left|Z_{x}^{(k)}(t)\right|^{2} d t=\sum_{m=1}^{N_{k}}\left|Y_{m}(x)\right|^{2}=\frac{N_{k}}{C_{d-1}} .
$$

Let $Y \in \mathcal{H}_{k}$. Then $Y(x)=\int_{S^{d-1}} Y(t) Z_{x}^{(k)}(t) d t$. By Schwarz's inequality,

$$
|Y(x)|^{2} \leq\left\|Z_{x}^{(k)}\right\|_{2}^{2}\|Y\|_{2}^{2} \leq \frac{N_{k}}{C_{d-1}}\|Y\|_{2}^{2}
$$

This establishes (iv).
สถาบันวิทยบริการ
จุฬาลงกรณ์มหาวิทยาลัย

## Chapter 4

## Laplacian on a sphere

We define the spherical Laplacian on $S^{d}$ to be an operator defined on $S^{d}$ given by the formula:

$$
\Delta_{S^{d}}=\sum_{k<l}\left(x_{k} \frac{\partial}{\partial x_{l}}-x_{l} \frac{\partial}{\partial x_{k}}\right)^{2}
$$

Notice that expressions like $\frac{\partial}{\partial x_{k}}$ do not make sense when applied to a function that is defined only on the sphere $S^{d}$. For a smooth function $f$ on $S^{d}$, extend $f$ smoothly to a neighborhood of $S^{d}$, then apply $x_{l} \frac{\partial}{\partial x_{k}}-x_{k} \frac{\partial}{\partial x_{l}}$, and then restrict again to $S^{d}$. Note that

$$
\left(x_{l} \frac{\partial}{\partial x_{k}}-x_{k} \frac{\partial}{\partial x_{l}}\right)|x|^{2}=2 x_{l} x_{k}-2 x_{k} x_{l}=0 .
$$

Thus these derivatives are all in directions tangentoto $S^{d}$. Therefore the values of the operator on $S^{d}$ is independent of the choice of extension.

We choose the domain of $\Delta_{S^{d}}$ to be the space $H$ of linear combinations of spherical harmonics on $S^{d}$. Since this space is dense in $L^{2}\left(S^{d}\right)$, it follows that $\Delta_{S^{d}}$ is a denselydefined operator on $L^{2}\left(S^{d}\right)$. Moreover, $\Delta_{S^{d}}$ is essentially self-adjoint, which we prove it later.

Proposition 4.1. Let $Y$ be a spherical harmonic of degree $k$. Then

$$
\Delta_{S^{d}} Y=-k(k+d-1) Y
$$

Proof.

$$
\begin{aligned}
-\Delta_{S^{d}} & =-\sum_{i<j}\left(x_{i} \frac{\partial}{\partial x_{j}}-x_{j} \frac{\partial}{\partial x_{i}}\right)^{2} \\
& =-\frac{1}{2} \sum_{i \neq j}\left(x_{i} \frac{\partial}{\partial x_{j}}-x_{j} \frac{\partial}{\partial x_{i}}\right)^{2} \\
& =-\frac{1}{2} \sum_{i \neq j}\left(x_{i}^{2} \frac{\partial^{2}}{\partial x_{j}^{2}}+x_{j}^{2} \frac{\partial^{2}}{\partial x_{i}^{2}}-x_{j}\left(\frac{\partial}{\partial x_{j}}+x_{i} \frac{\partial}{\partial x_{i}} \frac{\partial}{\partial x_{j}}\right)-x_{i}\left(\frac{\partial}{\partial x_{i}}+x_{j} \frac{\partial}{\partial x_{i}} \frac{\partial}{\partial x_{j}}\right)\right) \\
& =\frac{1}{2} \sum_{i \neq j}\left(-x_{i}^{2} \frac{\partial^{2}}{\partial x_{j}^{2}}-x_{j}^{2} \frac{\partial^{2}}{\partial x_{i}^{2}}+x_{j} \frac{\partial}{\partial x_{j}}+x_{i} \frac{\partial}{\partial x_{i}}+2 x_{i} x_{j} \frac{\partial}{\partial x_{i}} \frac{\partial}{\partial x_{j}}\right) \\
& =\sum_{i \neq j}-x_{i}^{2} \frac{\partial^{2}}{\partial x_{j}^{2}}+\sum_{i \neq j} x_{i} x_{j} \frac{\partial}{\partial x_{i}} \frac{\partial}{\partial x_{j}}+\frac{1}{2} \sum_{i \neq j}\left(x_{j} \frac{\partial}{\partial x_{j}}+x_{i} \frac{\partial}{\partial x_{i}}\right)
\end{aligned}
$$

For fixed $j$,

$$
\begin{gathered}
\sum_{\substack{i=1 \\
i \neq j)}}^{d+1}-x_{i}^{2} \frac{\partial^{2}}{\partial x_{j}^{2}}=-\left(x_{1}^{2}+x_{2}^{2}+\cdots+x_{d+1}^{2}\right) \frac{\partial^{2}}{\partial x_{j}^{2}}+x_{j}^{2} \frac{\partial^{2}}{\partial x_{j}^{2}} \\
\sum_{j=1}^{d+1} \sum_{\substack{i=1 \\
(i \neq j)}}^{d+1}-x_{i}^{2} \frac{\partial^{2}}{\partial x_{j}^{2}}=-|x|^{2} \Delta+\sum_{j=1}^{d+1} x_{j}^{2} \frac{\partial^{2}}{\partial x_{j}^{2}}, \quad \text { where } \Delta=\sum_{j=1}^{d+1} \frac{\partial^{2}}{\partial x_{j}^{2}} .
\end{gathered}
$$

Therefore

$$
-\Delta_{S^{d}}=-|x|^{2} \Delta+\sum_{j=1}^{d+1} x_{j}^{2} \frac{\partial^{2}}{\partial x_{j}^{2}}+\sum_{i \neq j} x_{i} x_{j} \frac{\partial}{\partial x_{i}} \frac{\partial}{\partial x_{j}}+\frac{1}{2} \sum_{i \neq j}\left(x_{j} \frac{\partial}{\partial x_{j}}+x_{i} \frac{\partial}{\partial x_{i}}\right)
$$



Since $Y$ is homogeneous, it is enough to prove the statement for only a harmonic monomial $P=x_{1}^{n_{1}} \ldots x_{d+1}^{n_{d+1}}$ where $n_{1}+\cdots+n_{d+1}=k$.

$$
\begin{aligned}
-\Delta_{S^{d}} P & =\left(\sum_{j=1}^{d+1} n_{j}\left(n_{j}-1\right)+2 \sum_{i<j} n_{i} n_{j}+\frac{1}{2} \sum_{i \neq j}\left(n_{i}+n_{j}\right)\right) P \\
& =\left(\sum_{j=1}^{d+1} n_{j}^{2}+2 \sum_{i<j} n_{i} n_{j}-\sum_{j=1}^{d+1} n_{j}+d \sum_{i=1}^{d+1} n_{i}\right) P \\
& =\left(\left(\sum_{j=1}^{d+1} n_{j}\right)^{2}-k+k d\right) P \\
& =\left(k^{2}-k+k d\right) P \\
& =k(k+d-1) P .
\end{aligned}
$$

To prove essential self-adjointness of $\Delta_{S^{d}}$, we recall the following theorem. For the definition of self-adjointness, a reader can refer to [6], p 256.

Theorem 4.2 ([6], p 257 ). Let $T$ be a densely defined symmetric operator on a Hilbert space $H$. Then the following statements are equivalent:
a $T$ is essentially self-adjoint;
b $\operatorname{Ker}\left(T^{*} \pm i\right)=0 ;$
c Range $(T \pm i)$ are dense in $H$. $9 /$ ? $9 \|$ ?
Now we are ready to verify the following proposition.
Proposition 4.3. The Laplacian $\Delta_{S^{d}}$ is essentially self-adjoint.

Proof. Firstly we verify that $\Delta_{S^{d}}$ is symmetric. Let

$$
X_{i j}=x_{i} \frac{\partial}{\partial x_{j}}-x_{j} \frac{\partial}{\partial x_{i}}
$$

Let $s o(d)$ be defined by $\left\{A \in M_{n}(\mathbb{R}) \mid A^{t}=-A\right\}$. From the definition of Lie derivative we can write

$$
X_{i j} f(x)=\left.\frac{d}{d t} f\left(e^{-t A_{i j}} x\right)\right|_{t=0}, \text { for some } A_{i j} \in s o(d+1) .
$$

If $f, g$ are smooth functions on $S^{d}$, then

$$
\begin{aligned}
\left\langle X_{i j} f, g\right\rangle & =\left.\int_{S^{d}}\left[\frac{d}{d t} f\left(e^{-t A_{i j}} y\right)\right]\right|_{t=0} g(y) d y \\
& =\left.\frac{d}{d t}\left[\int_{S^{d}} f\left(e^{-t A_{i j}} y\right) g(y) d y\right]\right|_{t=0} \quad \text { (by compactness of } S^{d} \text { ) } \\
& =\left.\frac{d}{d t}\left[\int_{S^{d}} f(y) g\left(e^{t A_{i j}} y\right) d y\right]\right|_{t=0} \quad \text { (since } d y \text { is rotationally invariant) } \\
& =-\left\langle f, X_{i j} g\right\rangle .
\end{aligned}
$$

In particular, $\left\langle\Delta_{S^{d}} f, g\right\rangle=\left\langle f, \Delta_{S^{d}} g\right\rangle$, for any spherical harmonics $f, g$. Thus $\Delta_{S^{d}}$ is symmetric on $H$. Now to show that $\Delta_{S^{d}}$ is essentially self-adjoint, it is enough to show that

$$
\overline{\text { Range }\left(\Delta_{S^{d}}+i\right)}=L^{2}\left(S^{d}, d x\right) .
$$

Recall that if $Y_{k}$ is a spherical harmonic of degree $k$, then

$$
\left(\Delta_{S^{d}}+i\right) Y_{k}=(-k(k+d-1)+i) Y_{k} .
$$

Let $\left\{Y_{1, k}, \ldots, Y_{N_{k}, k}\right\}$ be an orthonormal basis for $\mathcal{H}_{k}$. and $f \in L^{2}\left(S^{d}\right)$. We write

$$
\begin{array}{r}
f=\sum_{k=0}^{\infty} \sum_{n=1}^{N_{k}}\left\langle f, Y_{n, k}\right\rangle Y_{n, k} \\
=\sum_{k=0}^{\infty} \sum_{n=1}^{N_{k}}\left\langle f, Y_{n, k}\right\rangle \frac{\left(\Delta_{S^{d}}+i\right) Y_{n, k}}{(-k(k+d-1)+i)} .
\end{array}
$$

Hence $f \in \overline{\operatorname{Range}\left(\Delta_{S^{d}}+i\right)}$. This shows that $\Delta_{S^{d}}$ is essentially self-adjoint.

Now we turn to the Laplacian on a complex sphere. Let $J_{a}^{2}$ and $J_{\bar{a}}^{2}$ denote the differential operators on $S_{\mathbb{C}}^{d}$ given by

$$
\begin{aligned}
& J_{a}^{2}=\sum_{k<l}\left(a_{l} \frac{\partial}{\partial a_{k}}-a_{k} \frac{\partial}{\partial a_{l}}\right)^{2} \\
& J_{\bar{a}}^{2}=\sum_{k<l}\left(\bar{a}_{l} \frac{\partial}{\partial \bar{a}_{k}}-\bar{a}_{k} \frac{\partial}{\partial \bar{a}_{l}}\right)^{2} .
\end{aligned}
$$

In the same way as $\Delta_{S^{d}}, J_{a}^{2}$ and $J_{\bar{a}}^{2}$ can be interpreted as operators on $S_{\mathbb{C}}^{d}$.


## Chapter 5

## Segal-Bargmann transform

For each point $\mathbf{x} \in S^{d}$ and $t>0$, there is the heat kernel based at $\mathbf{x}$ denoted by $\rho_{t}(\mathbf{x}, \cdot)$ with the property that

$$
\begin{aligned}
\frac{d}{d t} \rho_{t}(\mathbf{x}, \mathbf{y}) & =\frac{1}{2} \Delta_{S^{d}} \rho_{t}(\mathbf{x}, \mathbf{y}), \quad \text { and } \\
\lim _{t \rightarrow 0^{+}} \int_{S^{d}} \rho_{t}(\mathbf{x}, \mathbf{y}) f(\mathbf{y}) d \mathbf{y} & =f(\mathbf{x}) \text { for any } f \in C\left(S^{d}\right)
\end{aligned}
$$

Given any function $f$ in $L^{2}\left(S^{d}\right)$ we define the Segal-Bargmann transform $C_{t} f$ of $f$ by

$$
\begin{equation*}
C_{t} f(\mathbf{a})=\int_{S^{d}} \rho_{t}(\mathbf{a}, \mathbf{x}) f(\mathbf{x}) d \mathbf{x}, \quad \mathbf{a} \in S_{\mathbb{C}}^{d} \tag{5.1}
\end{equation*}
$$

where $\rho_{t}$ is the heat kernel on $S^{d}$, with $\rho_{t}(\cdot, \mathbf{x})$ extended by analytic continuation from $S^{d}$ to $S_{\mathbb{C}}^{d}$.
It is easy to see that $C_{t} f$ is a holomorphic function on $S_{\mathbb{C}}^{d}$. Analogous to the SegalBargmann transform on $\mathbb{R}^{d}$, we expect that $C_{t}$ maps onto a space of holomorphic functions which are square-integrable with respect to a certain measure on $S_{\mathbb{C}}^{d}$. We can describe this measure explicitly using the identification $S_{\mathbb{C}}^{d} \cong T^{*}\left(S^{d}\right)$ via the
map

$$
\mathbf{a}(\mathbf{x}, \mathbf{p})=(\cosh p) \mathbf{x}+i \frac{\sinh p}{p} \mathbf{p} \quad \text { for any }(\mathbf{x}, \mathbf{p}) \in T^{*}\left(S^{d}\right) .
$$

Let $\nu_{t}$ be the function satisfying the following differential equation:

$$
\begin{equation*}
\frac{\partial}{\partial t} \nu_{t}(R)=\frac{1}{2}\left[\frac{\partial^{2}}{\partial R^{2}}+(d-1) \frac{\cosh R}{\sinh R} \frac{\partial}{\partial R}\right] \nu_{t}(R) \tag{5.2}
\end{equation*}
$$

with the initial condition

$$
\lim _{t \rightarrow 0^{+}} c_{d} \int_{0}^{\infty} f(R) \nu_{t}(R)(\sinh R)^{d-1} d R=f(0)
$$

for all continuous functions $f$ on $[0, \infty]$. Here $c_{d}$ is the volume of the unit sphere in $\mathbb{R}^{d}$. The existence of $\nu_{t}$ is guaranteeed by [2], Section 5.7. Then the desired measure can be written as

$$
\nu_{2 t}(2 p)\left(\frac{\sinh 2 p}{2 p}\right)^{d-1} 2^{d} d \mathbf{p} d \mathbf{x}
$$

where $d \mathbf{p}$ is Lebesgue measure on a d-dimensional real vecter space, $p=|\mathbf{p}|$ and $d \mathbf{x}$ is the surface measure on $S^{d}$.

In fact,

is the (complex) rotationally invariant measure on $S_{\mathbb{C}}^{d} \cong T^{*}\left(S^{d}\right)$ and $\nu_{2 t}(2 p)$ is the density with respect to this measure. We state it in the following proposition.

Proposition 5.1. The measure
is invariant under the action of $S O(d+1, \mathbb{C})$ on $S_{\mathbb{C}}^{d} \cong T^{*}\left(S^{d}\right)$.

To prove this proposition, we need to use the next lemma. Let $R=2 p$ and let $\alpha=|\mathbf{a}|^{2}=\sum\left|a_{k}\right|^{2}$. Then

$$
\alpha:=|\mathbf{a}|^{2}=\cosh ^{2} p+\sinh ^{2} p=\cosh 2 p .
$$

Since $p \geq 0$, we have $R=2 p=\cosh ^{-1} \alpha$. Hence $p$ can be considered as a function of $\mathbf{a}$.

Lemma 5.2. Let $\phi$ be a smooth, even, real-valued function on $\mathbb{R}$ and consider the function on $S_{\mathbb{C}}^{d}$ given by $p \mapsto \phi(2 p)$ where $p$ is regarded as a function of $\mathbf{a}$. Then

$$
J_{\mathbf{a}}^{2} \phi(2 p)=J_{\mathbf{a}}^{2} \phi(2 p)=-\left[\frac{\partial^{2} \phi}{\partial R^{2}}+(d-1) \frac{\cosh R}{\sinh R} \frac{\partial \phi}{\partial R}\right]_{R=2 p}
$$

Proof. By Chain Rule, we have

$$
\begin{aligned}
\frac{\partial \phi}{\partial a_{k}}(R) & =\frac{d \phi}{d R} \frac{d R}{d \alpha} \frac{\partial \alpha}{\partial a_{k}} \\
\frac{d R}{d \alpha} & =\frac{1}{\left(1-\alpha^{2}\right)^{\frac{1}{2}}}=-\frac{1}{\left(1-|\mathbf{a}|^{4}\right)^{\frac{1}{2}}} \\
\frac{\partial \alpha}{\partial a_{k}} & =\overline{a_{k}}
\end{aligned}
$$

Then

$$
\begin{gathered}
\left(a_{l} \frac{\partial}{\partial a_{k}}-a_{k} \frac{\partial}{\partial a_{l}}\right) \phi(R)=a_{l} \frac{d \phi}{d R}\left(-\frac{1}{\left(1-|\mathbf{a}|^{4}\right)^{\frac{1}{2}}}\right) \bar{a}_{k}-a_{k} \frac{d \phi}{d R}\left(-\frac{1}{\left(1-|\mathbf{a}|^{4}\right)^{\frac{1}{2}}}\right) \bar{a}_{l} \\
=\left(a_{k} \bar{a}_{l}=a_{l} \bar{a}_{k}\right) \frac{d \phi}{d R}\left(\frac{1}{\left(1\left|-|\mathbf{a}|^{4}\right)^{\frac{1}{2}}\right.}\right) .
\end{gathered}
$$

$$
\begin{aligned}
& \text { Consider } \\
& \frac{\partial}{\partial a_{k}}\left(a_{k} \overline{a_{l}}-a_{l} \bar{a}_{k}\right) \frac{d \phi}{d R}\left(\frac{9}{\left(1-|\mathbf{a}|^{4}\right)^{\frac{1}{2}}}\right) \\
& =\frac{d \phi}{d R}\left(\frac{-1}{\left(1-|\mathbf{a}|^{4}\right)^{\frac{1}{2}}}\right) \overline{a_{l}}+\left(a_{k} \bar{a}_{l}-a_{l} \bar{a}_{k}\right) \frac{d \phi}{d R} \frac{\partial}{\partial a_{k}}\left(\frac{1}{\left(1-|\mathbf{a}|^{4}\right)^{\frac{1}{2}}}\right)+\left(\frac{1}{\left(1-|\mathbf{a}|^{4}\right)^{\frac{1}{2}}}\right) \frac{\partial}{\partial a_{k}} \frac{d \phi}{d R} \\
& =\frac{d \phi}{d R}\left(\frac{1}{\left(1-|\mathbf{a}|^{4}\right)^{\frac{1}{2}}} \overline{a_{l}}\right)+\left(a_{k} \overline{a_{l}}-a_{l} \overline{a_{k}}\right)\left(|\mathbf{a}|^{2} \frac{d \phi}{d R}\left(\frac{\overline{a_{k}}}{\left(1-|\mathbf{a}|^{4}\right)^{\frac{3}{2}}}-\frac{\overline{a_{k}}}{\left(1-|\mathbf{a}|^{4}\right)^{\frac{3}{2}}} \frac{d^{2} \phi}{d R^{2}}\right) .\right.
\end{aligned}
$$

Then

$$
\begin{aligned}
& \left(a_{l} \frac{\partial}{\partial a_{k}}-a_{k} \frac{\partial}{\partial a_{l}}\right)^{2} \phi(R) \\
& =a_{l}\left(\frac{d \phi}{d R}\left(\frac{1}{\left(1-|\mathbf{a}|^{4}\right)^{\frac{1}{2}}} \overline{a_{l}}\right)+\left(a_{k} \overline{a_{l}}-a_{l} \overline{a_{k}}\right)\left(|\mathbf{a}|^{2} \frac{d \phi}{d R}\left(\frac{\overline{a_{k}}}{\left(1-|\mathbf{a}|^{4}\right)^{\frac{3}{2}}}-\frac{\overline{a_{k}}}{\left(1-|\mathbf{a}|^{4}\right)^{\frac{3}{2}}} \frac{d^{2} \phi}{d R^{2}}\right)\right)\right. \\
& +a_{k}\left(\frac{d \phi}{d R}\left(\frac{1}{\left(1-|\mathbf{a}|^{4}\right)^{\frac{1}{2}} \overline{a_{l}}}\right)+\left(a_{k} \overline{a_{l}}-a_{l} \overline{a_{k}}\right)\left(|\mathbf{a}|^{2} \frac{d \phi}{d R}\left(\frac{\overline{a_{k}}}{\left(1-|\mathbf{a}|^{4}\right)^{\frac{3}{2}}}-\frac{\overline{a_{k}}}{\left(1-|\mathbf{a}|^{4}\right)^{\frac{3}{2}}} \frac{d^{2} \phi}{d R^{2}}\right)\right)\right. \\
& =\left(\frac{a_{k} \overline{a_{l}}-a_{l} \overline{a_{k}}}{|\mathbf{a}|^{4}-1}\right) \frac{\partial^{2} \phi}{\partial R^{2}}-\frac{\left(\left|a_{k}\right|^{2}+\left|a_{l}\right|^{2}\right)\left(|\mathbf{a}|^{4}-1\right)+|\mathbf{a}|^{2}\left(a_{k} \overline{a_{l}}-a_{l} \overline{a_{k}}\right)^{2}}{\partial \phi} \frac{\partial \phi}{\partial R} .
\end{aligned}
$$

Note that

$$
\begin{aligned}
& \sum_{k<l}\left(\left|a_{k}\right|^{2}+\left|a_{l}\right|^{2}\right)=\frac{1}{2} \sum_{k \neq l}\left(\left|a_{k}\right|^{2}+\left|a_{l}\right|^{2}\right) \\
& =\frac{1}{2} \sum_{k, l}\left(1-\delta_{k l}\right)\left(\left|a_{k}\right|^{2}+\left|a_{l}\right|^{2}\right) \\
& =\frac{1}{2}\left[2(d+1)|\mathbf{a}|^{2}-|\mathbf{a}|^{2}\right] \\
& =d|\mathbf{a}|^{2} \text {. } \\
& \sum_{k<l}\left(a_{k} \overline{a_{l}}-a_{l} \overline{a_{k}}\right)^{2}=\frac{1}{2} \sum_{k, l}\left(1-\delta_{k l}\right)\left(a_{k}^{2}{\overline{a_{l}}}^{2}+a_{l}^{2}{\overline{a_{k}}}^{2}-2\left|a_{l}\right|^{2}\left|a_{k}\right|^{2}\right) \\
& =\frac{1}{2}\left(\left|a^{2}\right|^{2}+\left|a^{2}\right|^{2}-2|\mathbf{a}|^{4}\right)+\left(\sum_{k}\left|a_{k}\right|^{4}+\left|a_{k}\right|^{4}-2\left|a_{k}\right|^{4}\right) \\
& =-\left(|\mathbf{a}|^{4}-\left|a^{2}\right|^{2}\right)
\end{aligned}
$$

Then

$$
9 \sum_{k<l}\left(a_{l} \frac{6 \partial}{\partial a_{k}}-a_{k} \frac{\partial}{\partial a_{l}}\right)^{2} \phi(R)=-\frac{\partial^{2} \phi}{\partial R^{2}}-(d-1) \frac{\left.\square \mathbf{a}\right|^{2} \mid}{\sqrt{|\mathbf{a}|^{4}-1}} \frac{\partial \phi}{\partial R} .
$$

Since $|\mathbf{a}|^{2}=\cosh R, \sqrt{|\mathbf{a}|^{4}-1}=\sinh R$, so we obtain the claimed formula.
Let us recall the following theorem about Haar measure on a Lie group.

Theorem 5.3 ([5], Theorem 8.36). Let $G$ be a Lie group, let $H$ be a closed subgroup, and let $\triangle_{G}$ and $\triangle_{H}$ be the respective modular functions. Then a necessary and sufficient condition for $G / H$ to have a nonzero $G$-invariant Borel measure is that the restriction to $H$ of $\triangle_{G}$ is equal to $\triangle_{H}$. In this case such a measure $d \mu(g H)$ is unique up to a scalar multiplication.

Proof of Proposition 5.1. We consider $S_{\mathbb{C}}^{d}$ as the quotient $S O(d+1, \mathbb{C}) / S O(d, \mathbb{C})$. Since $S O(d+1, \mathbb{C})$ and $S O(d, \mathbb{C})$ are unimodular, it follows that the modular functions $\Delta_{S O(d+1, \mathrm{C})}$ and $\Delta_{S O(d, \mathrm{C})}$ are equal to 1. Hence by Theorem 5.3, there is a smooth $S O(d+1, \mathbb{C})$-invariant measure on $S_{\mathbb{C}}^{d}$ and it is unique up to a constant. Especially, this measure must be $S O(d+1)$-invariant. Since $d \mathbf{p} d \mathbf{x}$ is also $S O(d+1)$ invariant, this measure must be of the form $\gamma(p) d \mathbf{p} d \mathbf{x}$ for some smooth function $\gamma$. Let

and $g(p)=\frac{\alpha(p)}{\beta(p)}$. Therefore $g$ is an $S O(d+1)$-invariant function. We will consider $J_{\mathbf{a}}^{2}$ only on the space of $S O(d+1)$-invariant functions. We already know that $J_{\mathbf{a}}^{2}$ must be self-adjoint with respect to the $S O(d+1, \mathbb{C})$-invariant measure. In particular, $J_{\mathbf{a}}^{2}$ must be self-adjoint when restricted to the space of $S O(d+1)$-invariant functions, which we can write in the form $f(a)=\phi(2 p)$ for some smooth function $\phi$ on $\mathbb{R}$. By Lemma 5.2, on the space of $S O(d+1)$-invariant functions, $J_{\mathbf{a}}^{2}$ is equal to the hyperbolic Laplacian ([2], page 177). Since the measure $\beta(p) d \mathbf{p} d \mathbf{x}$ is just the hyperbolic volume measure, $J_{\mathbf{a}}^{2}$ is self-adjoint with respect to the measure $\beta(p) d \mathbf{p} d \mathbf{x}$. We can conclude that on $S O(d+1)$-invariant functions, $J_{\mathbf{a}}^{2}$ is self-adjoint with respect
to both the measure $\gamma(p) d \mathbf{p} d \mathbf{x}$ and $\beta(p) d \mathbf{p} d \mathbf{x}$. Then

$$
\begin{aligned}
\left\langle J_{a}^{2} g(p), J_{a}^{2} g(p)\right\rangle_{\beta} & =\int_{x \in S^{d}} \int_{p \cdot x=0} J_{a}^{2}(g(p)) J_{a}^{2}(g(p)) \beta(p) d \mathbf{p} d \mathbf{x} \\
& =\int_{x \in S^{d}} \int_{p \cdot x=0} J_{a}^{2}\left(J_{a}^{2}(g(p)) \frac{\alpha(p)}{\beta(p)} \beta(p) d \mathbf{p} d \mathbf{x}\right. \\
& =\int_{x \in S^{d}} \int_{p \cdot x=0} J_{a}^{2}\left(J_{a}^{2}(g(p)) \alpha(p) d \mathbf{p} d \mathbf{x}\right. \\
& =\int_{x \in S^{d}} \int_{p \cdot x=0} J_{a}^{2}(g(p)) J_{a}^{2}(1) \alpha(p) d \mathbf{p} d \mathbf{x}
\end{aligned}
$$

Thus $J_{a}^{2} g(p)=0$. By Lemma 5.2,

$$
\left.\left[\frac{\partial^{2} g}{\partial R^{2}}+(d-1) \frac{\cosh R}{\sinh R} \frac{\partial g}{\partial R}\right]\right|_{R=2 p}=0
$$

Since $g$ is a smooth $S O(d+1)$-invariant function on $S_{\mathbb{C}}^{d}, g$ is an even function. Then $\left.\frac{\partial g}{\partial R}\right|_{R=0}=0$. Solving the equation gives

$$
\frac{\partial g}{\partial R}=c e^{-(d-1) \int_{R}^{1} \operatorname{coth} S d S}
$$

Then

$$
\begin{aligned}
\left.\frac{\partial g}{\partial R}\right|_{R=0} & =c \lim _{\epsilon \rightarrow 0^{+}} e^{-(d-1) \int_{\epsilon}^{1} \operatorname{coth} S d S} \\
& =0
\end{aligned}
$$ Therefore $c=0$, and so is $\frac{\partial g}{\partial R}$. Hence $g$ is constant, which implies that $\gamma$ is a constant multiple of $\beta$.

Let $\mathcal{H} L^{2}\left(S_{\mathbb{C}}^{d}, \nu_{t}\right)$ be the space of holomorphic functions $F$ on $S_{\mathbb{C}}^{d} \cong T^{*}\left(S^{d}\right)$ for which

$$
\int_{x \in S^{d}} \int_{p \cdot x=0}|F(\mathbf{a}(\mathbf{x}, \mathbf{p}))|^{2} \nu_{2 t}(2 p)\left(\frac{\sinh 2 p}{2 p}\right)^{d-1} 2^{d} d \mathbf{p} d \mathbf{x}<\infty .
$$

Here is the main theorem of this work.

Theorem 5.4. The Segal-Bargmann transform $C_{t}$ defined by (5.1) is a unitary map from $L^{2}\left(S^{d}, d \mathbf{x}\right)$ onto $\mathcal{H} L^{2}\left(S_{\mathbb{C}}^{d}, \nu_{t}\right)$.

We divide the proof into two parts : isometry and surjectivity.

Proposition 5.5. $C_{t}$ is an isometry.
Proof. Since the map $\mathbf{a} \mapsto \rho_{t}(\mathbf{a}, \mathbf{x})$ is holomorphic for each $\mathbf{x} \in S^{d}$, we have

$$
J_{\overline{\mathbf{a}}}^{2} \rho_{t}(\mathbf{a}, \mathbf{x})=0 \text { for each } \mathbf{x} \in S^{d} .
$$

By definition of the heat kernel, we have

$$
\frac{1}{2} J_{\mathbf{a}}^{2} \rho_{t}(\mathbf{a}, \mathbf{x})=\frac{\partial}{\partial t} \rho_{t}(\mathbf{a}, \mathbf{x}) .
$$

Let $f, g \in L^{2}\left(S^{d}\right)$. Then

$$
\begin{gathered}
J_{\overline{\mathbf{a}}}^{2} C_{t} f(\mathbf{a})=\int_{S^{d}}\left(J_{\overline{\mathrm{a}}}^{2} \rho_{t}(\mathbf{a}, \mathbf{x})\right) f(\mathbf{x}) d \mathbf{x}=0 \\
\frac{1}{2} J_{\mathbf{a}}^{2} C_{t} f(\mathbf{a})=\int_{S^{d}}\left(\frac{1}{2} J_{\mathbf{a}}^{2} \rho_{t}(\mathbf{a}, \mathbf{x})\right) f(\mathbf{x}) d \mathbf{x}=\frac{\partial}{\partial t} C_{t} f(\mathbf{a}) .
\end{gathered}
$$

Therefore

$$
\begin{aligned}
& \frac{1}{2}\left(J_{\mathbf{a}}^{2}+J_{\mathbf{a}}^{2}\right) C_{t} f(\mathbf{a}) \overline{C_{t} g(\mathbf{a})}=\frac{1}{2}\left\{\left(J_{\mathbf{a}}^{2} C_{t} f(\mathbf{a})\right) \overline{C_{t} g(\mathbf{a})}+C_{t} f(\mathbf{a})\left(J_{\mathbf{a}}^{2} \overline{C_{t} g(\mathbf{a})}\right)+\right. \\
& \text { 6.6 } \underbrace{2}_{=} \frac{\partial}{\partial t} C_{t} f(\mathbf{a}) \frac{2}{C_{t}} C_{t} f(\mathbf{a})) \overline{C_{t} g(\mathbf{a})}+C_{t} f(\mathbf{a}) \frac{\partial}{\partial t} \overline{C_{t} g(\mathbf{a})}
\end{aligned}
$$

We know that $d \mathbf{a}=\beta(p) d \mathbf{p} d \mathbf{x}$ is $S O(d+1, \mathbb{C})$-invariant. Hence

$$
\begin{aligned}
& \frac{\partial}{\partial t} \int_{\mathbf{a} \in S_{\mathbb{C}}^{d}} C_{t} f(\mathbf{a}) \overline{C_{t} g(\mathbf{a})} \nu_{t}(\mathbf{a}) d \mathbf{a} \\
& =\int_{\mathbf{a} \in S_{\mathbb{C}}^{d}} \frac{\partial}{\partial t}\left(C_{t} f(\mathbf{a}) \overline{C_{t} g(\mathbf{a})} \nu_{t}(\mathbf{a})\right) d \mathbf{a} \\
& =\int_{\mathbf{a} \in S_{\mathrm{C}}^{d}} \frac{1}{2}\left(J_{\mathbf{a}}^{2}+J_{\overline{\mathbf{a}}}^{2}\right)\left(C_{t} f(\mathbf{a}) \overline{C_{t} g(\mathbf{a})}\right) \nu_{t}(\mathbf{a})+C_{t} f(\mathbf{a}) \overline{C g(\mathbf{a})} \frac{\partial \nu_{t}(\mathbf{a})}{\partial t} d \mathbf{a} \\
& =\int_{\mathbf{a} \in S_{\mathbb{C}}^{d}} C_{t} f(\mathbf{a}) \overline{C_{t} g(\mathbf{a})} \frac{1}{2}\left(J_{\mathbf{a}}^{2}+J_{\mathbf{a}}^{2}\right) \nu_{t}(\mathbf{a})+C_{t} f(\mathbf{a}) \overline{C_{t} g(\mathbf{a})} \frac{\partial \nu_{t}(\mathbf{a})}{\partial t} d \mathbf{a} \\
& =\int_{\mathbf{a} \in S_{\mathbb{C}}^{d}} C_{t} f(\mathbf{a}) \overline{C_{t} g(\mathbf{a})}\left(\frac{1}{2}\left(J_{\mathbf{a}}^{2}+J_{\overline{\mathbf{a}}}^{2}\right)+\frac{\partial}{\partial t}\right) \nu_{t}(\mathbf{a}) d \mathbf{a} \\
& =0 .
\end{aligned}
$$

From Lemma 5.2 and the differential equation (5.2) satisfied by $\nu_{t}$ we see that the last integral is zero. Then $\int_{\mathbf{a} \in S_{\mathrm{C}}^{d}} C_{t} f(\mathbf{a}) \overline{C_{t} g(\mathbf{a})} \nu_{t}(\mathbf{a}) d \mathbf{a}$ is independent of $t$. By the initial condition for $\nu$ and since $\rho_{t}(\mathbf{a}, \cdot)=\rho_{t}(\mathbf{x}, \cdot)$ when $p=0$ and $\lim _{t \rightarrow 0} \rho_{t}(\mathbf{x}, \cdot)=\delta_{\mathbf{x}}$,

$$
\lim _{t \rightarrow 0} \int_{\mathbf{a} \in S_{\mathbf{C}}^{d}} C_{t} f(\mathbf{a}) \overline{C_{t} g(\mathbf{a})} \nu_{t}(\mathbf{a}) d \mathbf{a}=\int_{S^{d}} f(\mathbf{x}) \overline{g(\mathbf{x})} d \mathbf{x} .
$$

Since the value of the first integral is independent of $t$, this shows that $C_{t}$ is isometric.

Next, we we turn to the proof of surjectivity of $C_{t}$. As we already know that if $Y_{k}$ is spherical harmonic of degree $k$,


$$
=\left(e^{\frac{t \lambda_{k}}{2}} Y_{k}\right)_{\mathbb{C}}
$$

$$
=e^{\frac{t \lambda_{k}}{2}}\left(Y_{k}\right)_{\mathbb{C}}, \text { where } \lambda_{k}=-k(k+d-1) .
$$

Thus the image of $C_{t}$ contains analytic continuation of all spherical harmonics. Since $C_{t}$ is an isometry, its image is a closed subspace of $L^{2}\left(S_{\mathbb{C}}^{d}, \nu_{t}\right)$. Thus it suffices to show that every holomorphic $L^{2}$ function on $S_{\mathbb{C}}^{d}$ can be approximated by spherical harmonics of holomorphic representations.

Let $F$ be any holomorphic function on $S_{\mathbb{C}}^{d}$. Then $\left.F\right|_{S^{d}}$ is a smooth function. $\left.F\right|_{S^{d}}$ can be written as a spherical harmonic expansion as follows:

$$
F(\mathrm{x})=\sum_{k=0}^{\infty} \sum_{n=1}^{N_{k}}\left\langle F, Y_{n, k}\right\rangle Y_{n, k}(\mathrm{x}) \quad\left(\mathrm{x} \in S^{d}\right)
$$

From Chapter 3 , we know that this series converges in $L^{2}\left(S^{d}, d \mathbf{x}\right)$. We will verify that this series converges uniformly.

Proposition 5.6. If $F$ is a smooth function on $S^{d}$, then the spherical harmonic expansion converges uniformly to $F$.

Proof.


This implies that 6169919 ² $\left\langle\Delta_{S^{d}}^{m} F, Y_{n, k}\right\rangle=(-1)^{m}(k(k+d-1))^{m}\left\langle F, Y_{n, k}\right\rangle$.

Then we can estimates the series by

$$
\begin{aligned}
& \sum_{k=0}^{\infty} \sum_{n=1}^{N_{k}}\left|\left\langle F, Y_{n, k}\right\rangle Y_{n, k}(\mathbf{x})\right| \\
& \leq \sum_{k=0}^{\infty} \sum_{n=1}^{N_{k}}\left|\left\langle F, Y_{n, k}\right\rangle\right| \sqrt{\frac{N_{k}}{C_{d}}}\left\|Y_{n, k}\right\|_{2} \\
& =\frac{1}{\sqrt{C_{d}}} \sum_{k=0}^{\infty} \sum_{n=1}^{N_{k}} \sqrt{N_{k}}\left|\left\langle F, Y_{n, k}\right\rangle\right| \\
& =\frac{1}{\sqrt{C_{d}}} \sum_{k=0}^{\infty} \sum_{n=1}^{N_{k}}\left(\frac{\sqrt{N_{k}}}{k^{d}(k+d-1)^{d}}\left|\left\langle\Delta_{S^{d}}^{d} F, Y_{n, k}\right\rangle\right|\right) \\
& \leq \frac{1}{\sqrt{C_{d}}}\left(\sum_{k=0}^{\infty} \frac{N_{k}^{2}}{k^{2 m}(k+d-1)^{2 m}}\right)^{\frac{1}{2}}\left(\sum_{k=0}^{\infty} \sum_{n=1}^{N_{k}}\left|\left\langle\Delta_{S^{d}}^{m} F, Y_{n, k}\right\rangle\right|^{2}\right)^{\frac{1}{2}} \\
& =\frac{1}{\sqrt{C_{d}}}\left(\sum_{k=0}^{\infty} \frac{N_{k}^{2}}{k^{2 m}(k+d-1)^{2 m}}\right)^{\frac{1}{2}}\left\|\Delta_{S^{d}}^{2 m} F\right\|_{2} .
\end{aligned}
$$

By (3.1), we know that

$$
N_{k}=\frac{\frac{(d+2 k-1)}{k}\binom{d+k-2}{k-1} . ~}{k}
$$

Hence,

$$
\begin{aligned}
& \sum_{k=0}^{\infty} \frac{N_{k}^{2}}{k^{2 m}(k+d-1)^{2 m}} \\
& =\frac{1}{((d-1)!)^{2}} \sum_{k=0}^{\infty} \frac{1}{k^{2 m}}\left(\frac{(k+d-2)!}{(k-1)!(k+d-1)^{m-2}}\right)^{2}\left(\frac{(d+2 k-1)}{k(d+k-1)^{2}}\right)^{2} .
\end{aligned}
$$

If $m \geq d$, then

$$
\frac{(k+d-2)!}{(k-1)!(k+d-1)^{m-2}}=\frac{(d+k-2) \ldots(k)}{(k+d-1)^{m-2}}<1 \text {. }
$$



$$
\frac{(d+2 k-1)}{k(d+k-1)^{2}} \rightarrow 0 \text { when } k \rightarrow \infty,
$$

the $k^{\text {th }}$ term in the series is bounded by $\frac{\text { const }}{k^{2 m}}$. By Comparison test, this series converges uniformly on $S^{d}$.

Since $F$ is assumed holomorphic, and each term in the Fourier series above has an analytic continuation to $S_{\mathbb{C}}^{d}$, it is natural to suggest the following series expansion for $F(a)$

$$
F(a)=\sum_{k=0}^{\infty} \sum_{n=1}^{N_{k}}\left\langle F, Y_{n, k}\right\rangle Y_{n, k}(a)
$$

We will refer to this series as the holomorphic Fourier series of $F$, abbreviated HFS.

Proposition 5.7. For any holomorphic function $F$ on $S_{\mathbb{C}}^{d}$ the holomorphic Fourier series of $F$ converges to $F$ uniformly on compact sets.

Proof. For each $a \in S_{\mathbb{C}}^{d}$, define the function $F_{a}(\mathbf{x})=F(A \mathbf{x})$, where $a=A e_{d+1} \cong$ $A S O(d, \mathbb{C}) . F_{a}$ is a smooth function on $S^{d}$. We will consider the Fourier series of $F_{a}$ at point $\mathbf{x}$ and HFS of $F$ at point $A \mathbf{x}$. Since Laplacian on $\mathbb{R}^{d}$ is invariant under the action of orthogonal matrix, $Y_{n, k}(A \mathbf{x})$ is a spherical harmonic of the same degree. Then

$$
Y_{n, k}(A \mathrm{x})=\sum_{l=1}^{N_{k}} a_{n, l}^{k} Y_{l, k}(\mathbf{x}),
$$

and

$$
\begin{aligned}
\sum_{l=1}^{N_{k}} a_{n, l}^{k} a_{m, l}^{k} & =\int_{S^{d}} Y_{n, k}(A \mathbf{x}) Y_{m, k}(A \mathbf{x}) d \mathbf{x} \\
& =\int_{S^{d}} Y_{n, k}(\mathbf{x}) Y_{m, k}(\mathbf{x}) d \mathbf{x} \\
6 & =\delta_{m n} .
\end{aligned}
$$

Thus $\left(a_{n, l}^{k}\right)_{1 \leq n, l \leq N_{k}}$ is an orthogonal matrix. We can write $Y_{n, k}(\mathbf{x})$ in the form

$$
Y_{n, k}(\mathbf{x})=\sum_{l=1}^{N_{k}} a_{n, l}^{k} Y_{l, k}\left(A^{-1} \mathbf{x}\right)
$$

Since $\left(a_{n, l}^{k}\right)^{-1}=\left(a_{n, l}^{k}\right)^{t}, Y_{n, k}\left(A^{-1} \mathbf{x}\right)=\sum_{l=1}^{N_{k}} a_{l, n}^{k} Y_{l, k}(\mathbf{x})$. Consider the coefficients of
the Fourier series of $F_{a}$ at point $\mathbf{x}$. We have

$$
\begin{aligned}
\left\langle F_{a}, Y_{n, k}\right\rangle & =\int_{S^{d}} F(A \mathbf{x}) Y_{n, k}(\mathbf{x}) d \mathbf{x} \\
& =\int_{S^{d}} F(\mathbf{x}) Y_{n, k}\left(A^{-1} \mathbf{x}\right) d \mathbf{x} \\
& =\sum_{l=1}^{N_{k}} a_{l, n}^{k} \int_{S^{d}} F(\mathbf{x}) Y_{l, k}(\mathbf{x}) d \mathbf{x} .
\end{aligned}
$$

This implies that

$$
\begin{aligned}
\sum_{k=0}^{\infty} \sum_{n=1}^{N_{k}}\left\langle F_{a}, Y_{n, k}\right\rangle Y_{n, k}(\mathrm{x}) & =\sum_{k=0}^{\infty} \sum_{n=1}^{N_{k}}\left(\sum_{l=1}^{N_{k}}\left\langle F, Y_{l, k}\right\rangle Y_{n, k}(\mathbf{x})\right) \\
& =\sum_{k=0}^{\infty} \sum_{l=1}^{N_{k}}\left\langle F, Y_{l, k}\right\rangle Y_{l, k}(A \mathbf{x}) .
\end{aligned}
$$

We can see that the holomorphic Fourier series of $F$ at the point $A \mathbf{x}$ is the same as the ordinary Fourier series of $F_{a}$ at the point $\mathbf{x}$.

Let $\mathbf{x}=e_{d+1}$. Then

$$
F(a)=F_{a}\left(e_{d+1}\right)=\text { HFS of } F \text { at } A e_{d+1}=\operatorname{HFS} \text { of } F \text { at } a .
$$

This shows that the HFS converges pointwise. We will verify that HFS converges uniformly on compact sets in $S_{\mathbb{C}}^{d}$. Let $K$ be any compact set in $S_{\mathbb{C}}^{d}$. Since $K$ is compact, there exists $M>0$ such that $\left|\Delta^{m} F(a)\right| \leq M$ for all $a \in K$. Because HFS of $F$ at $a$ is equal to the Fourier series of $F_{a}$ at $e_{d+1}$, it suffices to show that the series $\sum_{k=0}^{\infty} \sum_{m=1}^{N_{k}}\left\langle F_{a}, Y_{n, k}\right\rangle Y_{n, k}\left(e_{d+1}\right)$ as a function of $a$, converges uniformly on
$K$.

$$
\begin{aligned}
& \sum_{k=0}^{\infty} \sum_{n=1}^{N_{k}}\left|\left\langle F_{a}, Y_{n, k}\right\rangle Y_{n, k}\left(e_{d+1}\right)\right| \\
& \leq \sum_{k=0}^{\infty} \sum_{n=1}^{N_{k}}\left|\left\langle F_{a}, Y_{n, k}\right\rangle\right| \sqrt{\frac{N_{k}}{C_{d}}}\left\|Y_{n, k}\right\|_{2} \\
& =\frac{1}{\sqrt{C_{d}}} \sum_{k=0}^{\infty} \sum_{n=1}^{N_{k}} \frac{\sqrt{N_{k}}}{k^{m}(k+d-1)^{m}}\left|\left\langle\Delta^{m} F_{a}, Y_{n, k}\right\rangle\right| \\
& \leq \frac{1}{\sqrt{C_{d}}}\left(\sum_{k=0}^{\infty} \sum_{n=1}^{N_{k}}\left(\frac{N_{k}}{k^{2 m}(k+d-1)^{2 m}}\right)\right)^{\frac{1}{2}}\left(\sum_{k=0}^{\infty} \sum_{n=1}^{N_{k}}\left|\left\langle\Delta^{m} F_{a}, Y_{n, k}\right\rangle\right|^{2}\right)^{\frac{1}{2}} \\
& =\frac{1}{\sqrt{C_{d}}}\left(\sum_{k=0}^{\infty} \frac{N_{k}^{2}}{\left.k^{2 m}(k+d-1)^{2 m}\right)^{\frac{1}{2}}\left\|\Delta^{m} F_{a}\right\|_{2}}\right. \\
& \leq \frac{1}{\sqrt{C_{d}}}\left(\sum_{k=0}^{\infty} \frac{N_{k}^{2}}{\left.k^{2 m}(k+d-1)^{2 m}\right)^{\frac{1}{2}}\left\|\Delta^{m} F_{a}\right\|_{\infty}}\right. \\
& \leq \frac{1}{\sqrt{C_{d}}}\left(\sum_{k=0}^{\infty} \frac{N_{k}^{2}}{\left.k^{2 m}(k+d-1)^{2 m}\right)^{\frac{1}{2}} M .}\right.
\end{aligned}
$$

By proving in a similar way, we have holomorphic Fourier series of $F$ converges uniformly on any compact subset $K$ of $S_{\mathbb{C}}^{d}$.

Proposition 5.8. If $F \in \mathcal{H} L^{2}\left(S_{\mathbb{C}}^{d}, \nu_{t}\right)$, then the holomorphic Fourier series of $F$ converges to $F$ in $L^{2}\left(S_{\mathbb{C}}^{d}, \nu_{t}\right)$.

Proof. At first, we claim that this series is an orthogonal series. Let $Y_{k}$ and $Y_{l}$ be spherical harmonic polynomials of degree $k$ and $l$ respectively, $k \neq l$. Since $v_{t}(p) d a$ is $S O(d+1)$-invariant,

$$
\begin{aligned}
& =\left\langle\Delta_{S^{d}} Y_{k}, Y_{l}\right\rangle \\
& =\int_{S_{\mathrm{C}}^{d}} \Delta_{S^{d}} Y_{k}(a) \overline{Y_{l}(a)} v_{t}(a) d a .
\end{aligned}
$$

Since $\Delta_{S^{d}}$ commutes with analytic continuation,

$$
\Delta_{S^{d}}\left(Y_{n, k}\right)_{\mathbb{C}}=\left(\Delta_{S^{d}} Y_{n, k}\right)_{\mathbb{C}}=-k(k+d-1)\left(Y_{n, k}\right)_{\mathbb{C}}
$$

Then

$$
-l(l+d-1) \int_{S_{\mathbb{C}}^{d}} Y_{k}(a) \overline{Y_{l}(a)} v_{t}(a) d a=-k(k+d-1) \int_{S_{\mathbb{C}}^{d}} Y_{k}(a) \overline{Y_{l}(a)} v_{t}(a) d a .
$$

Since $l \neq k, \int_{S_{\mathbb{C}}^{d}} Y_{k}(a) \overline{Y_{l}(a)} v_{t}(a) d a=0$. Since the series is orthogonal, it will converge provided that the sum of the squares of the norms is finite.

Let $E_{n}=\left\{\mathbf{a}(\mathbf{x}, \mathbf{p}) \in S_{\mathbb{C}}^{d}| | \mathbf{p} \mid \leq n\right\}$. Claim that $E_{n}$ is $S O(d+1)$-invariant. Let $A \in S O(d+1)$. Then

$$
\begin{aligned}
A(\mathbf{a}(\mathbf{x}, \mathbf{p})) & =A \cosh p \mathbf{x}+i A \frac{\sinh p}{p} \mathbf{p} \\
& =\cosh p(A \mathbf{x})+i \frac{\sinh p}{p}(A \mathbf{p})
\end{aligned}
$$

Since $A \in S O(d+1)$, it follows that $|A \mathbf{p}|=|\mathbf{p}| \leq n,|A \mathbf{x}|=|\mathbf{x}|=1$ and $A \mathbf{x} \cdot A \mathbf{p}=$ $\mathbf{x} \cdot \mathbf{p}=0$. Thus $A(\mathbf{a}(\mathbf{x}, \mathbf{p}))=\mathbf{a}(A \mathbf{x}, A \mathbf{p}) \in E_{n}$.

Hence $E_{n}$ is an increasing sequence of compact $S O(d+1)$-invariant sets, with $\cup_{n} E_{n}=$ $S_{\mathbb{C}}^{d}$. Since $F=\sum_{k=0}^{\infty} Y_{k}$ on $E_{n}$, where $Y_{k}=\sum_{n=1}^{N_{k}}\left\langle F, Y_{n, k}\right\rangle Y_{n, k}$,

$$
\left\|\left.F\right|_{E_{n}}\right\|^{2}=\sum_{k=0}^{\infty}\left\|\left.Y_{k}\right|_{E_{n}}\right\|^{2}
$$

Note that $\left\|\left.Y_{k}\right|_{E_{n}}\right\|_{0}^{2}$ increases with $n$. 9 c/
We may apply Monotone Convergence Theorem so that

$$
\text { 99ำ } 9 \text { ¢ } \sum_{k=0}^{\infty} \psi Y_{k}\left\|^{2}=\right\| F \|^{2}<\infty
$$

Hence this series converges in $L^{2}\left(S_{\mathbb{C}}^{d}, \nu_{t}\right)$.

This completes the proof of surjectivity.

## Bibliography

[1] Axler, S., Bourdon, P. and Ramey, W. Harmonic Function Theory; New York, Springer-Verlag: 2001.
[2] Davies, E.B. Heat Kernels and Spectral Theory; Cambridge Cambridge, University Press: 1994.
[3] Hall, B. The Segal-Bargmann "coherent state" transform for compact Lie groups, J. Funct. Anal. 122: 1994, 103-151. MR 95e:22020
[4] Hall, B. and Mitchell, J. Coherent states on spheres, J. Math. Phys. 43: 2002, 1211-1236.
[5] Knapp, A.W. Lie Groups Beyond an Introduction; Boston, Birkhäuser: 1996.
[6] Reed, M. and Simon, B. Methods of Modern Mathematical Physics I: Functional Analysis; San Diego, Academic Press: 1972.
[7] Stenzel, M. The Segal-Bargmann transform on a symmetric space of compact type, J. Funct. Anal. 165: 1999, 44-58.
[8] Stein, E.M. and-Weiss, G. Introduction to Fourier-Analysis on Euclidean Spaces; Princeton, Princeton University Press: d971.
[9] Warner, F.W. Foundations of Differentiable Manifolds and Lie Groups; New York, Springer-Verlag: 1994.6 90 9

## VITA

Mr Boonyong Sriponpaew was born on May 6, 1980 in Bangkok, Thailand. He got a Bachelor of Science in Mathematics from Chulalongkorn University in 2001, and then he furthered his study for the Master's degree at the same place. Throughout his study, both undergraduate and graduate, he got the financial support from the Development and Promotion of Science and Talents Project (DSPT).


## จุฬาลงกรณ์มหาวิทยาลัย

