CHAPTER III

THE SCATTERING OF A PLANE WAVE FROM A CYLINDER

Suppose we have a plane wave which is incident on a circular object and is scattered by the object without change in frequency. The complete wave, satisfying the requirement that at large distances from the object it consists of a plane wave in the positive x-direction plus an outgoing radial wave or scattered wave, and that its normal gradient is zero at $r = a$, has the general form

where $\mathcal{U}_{\text{inc}}(r, \emptyset)$ is the incident wave,

 \mathcal{H}_{κ} (r, \emptyset) is the time independent scattered wave which satisfies the Helmholtz's equation,

and A is the amplitude of the incident wave.

Carrying through the straightforward separation of variables we obtain the following differential equations for the time independent scattered wave $R(r)$ $\oint (\emptyset)$

$$
\frac{d^{2}R(r)}{dr^{2}} + \frac{1}{r} \frac{dR(r)}{dr} + \frac{(k^{2} - n^{2}) R(r)}{r} = 0 \dots \dots \dots \dots 2)
$$

and

By the symmetry of the waves we shall require only solutions of 3) of the form

$$
\Phi(\emptyset) = B_n \cos n \emptyset \qquad \qquad \ldots \qquad \qquad \ldots \qquad \qquad \downarrow)
$$

where n is an integer to ensure that $\vec{\Phi}(2\pi - \vec{p}) = \vec{\Phi}(\vec{p})$.

In order to obtain solutions that represent circular waves we must use the following general solution of 2)

$$
R(r) = C_{n}H_{n}^{(1)}(kr) + D_{n}H_{n}^{(2)}(kr) \qquad \ldots \ldots \ldots \qquad (5)
$$

where $\pi_n^{(1)}(kr)$ and $\pi_n^{(2)}(kr)$ are the Hankel functions of the first and second kind respectively. We find by using the asymptotic form of the Hankel functions as r approaches infinity that at a large distance from the scattering circle R(r) e^{-iot} approaches

$$
C_{n} \left(\frac{2}{\pi k r}\right)^{1/2} e^{i(kr - \frac{n\pi}{2} - \frac{\pi}{4} - \omega t)} + D_{n} \left(\frac{2}{\pi k r}\right)^{1/2} e^{-i(kr - \frac{n\pi}{2} - \frac{\pi}{4} + \omega t)}
$$

Since the solution must represent outgoing waves*^{*} only we put $D_n = 0$. Thus we obtain the following time independent scattered wave function

$$
Y_s(r,\emptyset) = \sum_{n=0}^{\infty} B_n f_n^{(1)}
$$
 (kr) cos n \emptyset (6)

by putting $C_n = 1$ in equation 5). Therefore the complete time independent wave is

$$
u \quad u = \sum_{n=0}^{\infty} \left[A \epsilon_n i^n \, J_n(kr) + B_n H_n^{(1)}(kr) \, J \cos n\emptyset \quad \ldots \ldots \qquad 7 \right)
$$

where $\epsilon_n = 1$ when $n = 0$

 $= 2$ when $n > 0$.

Morse Philip M., Feshbach Herman. Methods of Theoretical Physics Part II (McGraw-Hill Book Company, Inc. 1953) pp.1377.

Consider the derivative of u

$$
\frac{\partial u}{\partial r} = \sum_{\substack{n=0 \ n=0}}^{\infty} \left[\underline{A} \xi_n i^n \frac{d}{dr} J_n(kr) + B_n \frac{d}{dr} J_n^{(1)}(kr) \right] \cos n \phi
$$

$$
= \sum_{n=0}^{\infty} \left[\underline{A} \xi_n i^n \frac{1}{2} \left\{ J_{n-1}(kr) - J_{n+1}(kr) \right\} +
$$

$$
\frac{1}{2} J_n \left\{ J_{n-1}^{(1)}(kr) - J_{n+1}^{(1)}(kr) \right\} \right]
$$

We shall assume a physical boundary condition of the form $\frac{\partial u}{\partial r}$ = 0 at $r = a$; therefore

$$
\frac{1}{2} A \xi_n^{i^n} \overline{[J_{n-1} (ka) - J_{n+1} (ka)]} + \frac{1}{2} B_n \overline{[J_{n-1} (ka) - J_{n+1} (ka)]} = 0
$$

that is $B_n = \frac{A \xi_n^{i^{n-2}} \overline{[J_{n-1} (ka) - J_{n+1} (ka)]}}{J_{n-1} (ka) - J_{n+1} (ka)}$ (8)

Therefore the complete wave thas the general form

where
$$
u_{\rm s} = A \sum_{n=0}^{\infty} \epsilon_n i^{n-2} \frac{\left[\frac{J}{J_{n-1}(ka)} - J_{n+1}(ka)\right]_n^{\rm (k)}}{\left[\frac{H}{J_{n-1}(ka)} - H_{n+1}(ka)\right]} \exp(-\frac{I_{n+1}(ka)}{I_{n+1}(ka)})
$$

Numerical Computation of scattered waves

We shall illustrate the calculation of the scattered wave by using a partial sum of \mathcal{U}_{s} .

Letting $k = 1$, $a = 6$, we have from 10)

$$
\mathbf{H}_{\rm s} = \frac{\sum_{n=0}^{\infty} A \in \mathbf{n}^{1-2} \left[J_{n-1}(6) - J_{n+1}(6) \right]}{\mathbf{H}_{n-1}^{(1)}(6) - \mathbf{H}_{n+1}^{(1)}(6)} \cdot \mathbf{H}_{n}^{(1)}(\mathbf{r}) \cos n\emptyset \dots 11}
$$

Without loss of generality we may assume the amplitude of the incident wave to be unity. Since

and
$$
J_{-n}(r) = (-1)^n J_n(r) , \tilde{Y}_{-n}^{\dagger}(r) = (-1)^n Y_n(r) \dots 13)
$$

we obtain the following formula for the m-th partial sum of 11)

$$
(7f_s)_m = \sum_{n=0}^{m!} \underbrace{f_n i^{n-2} \underbrace{J_{n-1}(6) - J_{n+1}(6)} \underbrace{J_n(r) + iY_n(r)}_{n-1}(6) - Y_{n+1}(6)}_{\text{no-1}}.
$$

 $m = 8$. We have Let

$$
(\mathcal{U}_{s})_{8} = \frac{\boxed{1^{-2} \quad J_{-1}(6) - J_{1}(\overline{6}) \boxed{J_{0}}(r) + i \quad Y_{0}(r)}}{\boxed{J_{-1}(6) - J_{1}(\overline{6}) + \boxed{Y_{-1}(6) - Y_{1}(\overline{6})}}
$$

$$
2 \t i^{-1} \left[J_0(6) - J_2(6) \right] J_1(r) + i Y_1(r) \cos \beta
$$

$$
J_0(6) - J_2(6) + i Y_0(6) - Y_2(6)
$$

$$
+\frac{2i^{o} \left[\overline{J}_{1}(6) - \overline{J}_{3}(6)\right] \left[\overline{J}_{2}(r) + i \overline{Y}_{2}(r)\right] \cos 2 \beta}{\left[\overline{J}_{1}(6) - \overline{J}_{3}(6)\right] + i \left[\overline{Y}_{1}(6) - \overline{Y}_{3}(6)\right]}
$$

$$
+\frac{2i^{1} \left[\overline{J}_{2}(6) - \overline{J}_{4}(6)\right] \left[\overline{J}_{3}(r) + i \overline{Y}_{3}(r)\right] \cos 3 \beta}{\left[\overline{J}_{2}(6) - \overline{J}_{4}(6)\right] + i \left[\overline{Y}_{2}(6) - \overline{Y}_{4}(6)\right]}
$$

$$
+\frac{2i^{2}\left[J_{3}(6)-J_{5}(6)\left[J_{4}(r)+i Y_{4}(\frac{r}{2})\right]\cos 4\beta}{\left[J_{2}(6)-J_{5}(6)+i\left[Y_{4}(6)-Y_{5}(6)\right]\right]} + \frac{2i^{3}\left[J_{4}(6)-J_{6}(6)\right]\left[J_{5}(r)+i Y_{5}(\frac{r}{2})\right]\cos 5\beta}{\left[J_{4}(6)-J_{6}(6)+i\left[Y_{4}(6)-Y_{6}(6)\right]\right]} + \frac{2i^{4}\left[J_{5}(6)-J_{7}(6)\right]\left[J_{6}(r)+i Y_{6}(\frac{r}{2})\right]\cos 6\beta}{\left[J_{5}(6)-J_{7}(6)\right]+i\left[Y_{5}(6)-Y_{7}(6)\right]} + \frac{2i^{5}\left[J_{6}(6)-J_{8}(6)\right]\left[J_{7}(r)+i Y_{7}(\frac{r}{2})\right]\cos 7\beta}{\left[J_{6}(6)-J_{8}(6)\right]+i\left[Y_{6}(6)-Y_{8}(6)\right]} + \frac{2i^{6}\left[J_{7}(6)-J_{9}(6)\right]\left[J_{8}(r)+i Y_{8}(\frac{r}{2})\right]\cos 8\beta}{\left[J_{7}(6)-J_{9}(6)\right]+i\left[Y_{7}(6)-Y_{9}(6)\right]} + \frac{2i^{6}\left[J_{7}(6)-J_{9}(6)\right]\left[J_{8}(r)+i Y_{8}(\frac{r}{2})\right]\cos 8\beta}{\left[J_{2}(6)-J_{9}(6)\right]+i\left[Y_{7}(6)-Y_{9}(6)\right]} + \frac{2(-0.2767)\left[J_{0}(r)+i Y_{0}(\frac{r}{2})\right]\left[J_{1}(r)+i Y_{1}(\frac{r}{2})\right]\cos\beta}{\left[0.1506-(-0.2429)\right]+i\left[-0.2882-0.2299\right]} + \frac{2[-0.2767-0.1148]\left[J_{2}(r)+i Y_{2}(\frac{r}{2})\right]\cos 2\beta}{\left[-0.2767-0.1148\right]+i\left[-0.1750-0.3283\right]} + \frac{2i\left[-0.2429-0.3576\right]+\frac{i\left[0.2299-0.0984\right]}{0.1148-0.3621\right]+i
$$

$$
\frac{2i\left[0.3576 - 0.2458\right] \left[\frac{1}{5}(r) + i \frac{r}{5}(r) \right] \cos 5\varnothing}{\left[0.3576 - 0.2458\right] + i\left[0.0984 - (-0.4268)\right]}
$$
\n
$$
2\left[0.3621 - 0.1296\right] \left[\frac{1}{5}(r) + i \frac{r}{6(r)} \right] \cos 6\varnothing
$$
\n
$$
+\frac{2i\left[0.3621 - 0.1296\right] \left[\frac{1}{5}(r) + i \frac{r}{6(r)} \right] \cos 6\varnothing
$$
\n
$$
+\frac{2i\left[0.2458 - 0.0565\right] \left[\frac{1}{5}(r) + i \frac{r}{7}(r) \right] \cos 7\varnothing}{\left[0.2458 - 0.0565\right] + i\left[0.4268 - (-1.1052)\right]}
$$
\n
$$
-\frac{2\left[0.1296 - 0.0212\right] \left[\frac{1}{3}(r) + i \frac{r}{8}(r) \right] \cos 8\varnothing}{\left[0.1296 - 0.0212\right] + i\left[0.6566 - (-2.2907)\right]}
$$
\n
$$
-\frac{0.35534}{\left[0.5554 + 0.35001\right]} \left[\frac{1}{0}(r) + i \frac{r}{1}(r) \right] \cos \varnothing
$$
\n
$$
-\frac{0.7870 i}{\left[0.3935 - 0.5181 i}\right] \left[\frac{1}{1}(r) + i \frac{r}{1}(r) \right] \cos \varnothing
$$
\n
$$
-\frac{0.7829}{-0.3914 - 0.5033 i}\left[\frac{1}{3}(r) + i \frac{r}{2}(r) \right] \cos 2\varnothing
$$
\n
$$
+\frac{0.4946}{-0.6005 + 0.1315 i}\left[\frac{1}{3}(r) + i \frac{r}{3}(r) \right] \cos 3\varnothing
$$
\n
$$
+\frac{0.4946}{\left[0.22
$$

$$
+\frac{0.3786 \text{ i}}{0.1893 + 0.6784 \text{ i}} \left[\frac{1}{J_7}(\mathbf{r}) + i \mathbf{Y}_7(\mathbf{r})\right] \cos 7 \beta
$$

\n
$$
-\frac{0.2168}{0.1084 + 1.6341 \text{ i}} \left[\frac{1}{3}(5) + i \mathbf{Y}_8(\mathbf{r})\right] \cos 8 \beta
$$

\n
$$
-\frac{(0.5534 \times 0.5534) - (0.5534 \times 0.3500 \text{ i}) \cdot [\mathbf{J}_0(\mathbf{r}) + i \mathbf{Y}_0(\mathbf{r})]}{(0.5534)^2 + (0.3500)^2}
$$

\n
$$
-\frac{(0.7870 \text{ i} \times 0.3935) + (0.7870 \text{ i} \times 0.5181 \text{ i}) [\mathbf{J}_1(\mathbf{r}) + i \mathbf{Y}_1(\mathbf{r})] \cos \beta}{(0.3935)^2 + (0.5181)^2}
$$

\n
$$
+\frac{(0.7829 \times 0.3914) - (0.7829 \times 0.50331) [\mathbf{J}_2(\mathbf{r}) + i \mathbf{Y}_2(\mathbf{r})] \cos 2\beta}{(0.3914)^2 + (0.5181)^2}
$$

\n
$$
+\frac{(1.2010 \times 0.6005) + (1.2010 \times 0.13151) [\mathbf{J}_3(\mathbf{r}) + i \mathbf{Y}_3(\mathbf{r})] \cos 3\beta
$$

\n
$$
-(0.4946 \times 0.2473) + (0.4946 \times 0.52541) [\mathbf{J}_4(\mathbf{r}) + i \mathbf{Y}_4(\mathbf{r})] \cos 4 \beta
$$

\n
$$
-(0.2473)^2 + (0.5254)^2
$$

\n
$$
-\frac{(0.2493)^2 + (0.5254)^2}{(0.1118)^2 + (0.5252)^2}
$$

\n
$$
+\frac{(0.4650 \times 0.2355) -
$$

$$
-\frac{0.3065 - 0.1937 \text{ i}}{0.3065 + 0.1225}
$$
\n
$$
+\frac{0.4077 - 0.3097 \text{ i}}{0.1548 + 0.2684}
$$
\n
$$
-\frac{0.4077 - 0.3097 \text{ i}}{0.1548 + 0.2684}
$$
\n
$$
-\frac{0.3064 - 0.3940 \text{ i}}{0.1532 + 0.2533}
$$
\n
$$
+\frac{0.7212 + 0.1579 \text{ i}}{0.3606 + 0.0173}
$$
\n
$$
-\frac{0.1223 + 0.2599 \text{ i}}{0.0612 + 0.2760}
$$
\n
$$
-\frac{0.1174 + 0.0250 \text{ i}}{0.0125 + 0.2758}
$$
\n
$$
+\frac{0.1081 - 0.2137 \text{ i}}{0.0541 + 0.2111}
$$
\n
$$
-\frac{0.2568 + 0.0717 \text{ i}}{0.0358 + 0.4602}
$$
\n
$$
-\frac{0.0235 - 0.3543 \text{ i}}{0.0358 + 0.4602}
$$
\n
$$
-\frac{0.0235 - 0.3543 \text{ i}}{0.018 + 0.4602}
$$
\n
$$
-\frac{0.0235 - 0.3543 \text{ i}}{0.0118 + 0.4602}
$$
\n
$$
-\frac{0.0235 - 0.3543 \text{ i}}{0.0118 + 0.4602}
$$
\n
$$
-\frac{0.0235 - 0.3543 \text{ i}}{0.0118 + 0.4602}
$$
\n
$$
-\frac{0.0235 - 0.3543 \text{ i}}{0.0118 + 0.4602}
$$

+ $(0.9634 - 0.7318 i)(J_1(r) + i Y_1(r))$ cos Ø + $(0.7538 - 0.9692 i)(J_2(r) + i Y_2(r) \cos 2 \emptyset)$ + $(1.9084 + 0.4178 i)(J_3(r) + i Y_3(r) \cos 3 \emptyset$

=

+ $(-0.3627 - 0.7708 i)(J_{\phi}(r) + i Y_{\mu}(r))$ cos 4p + $(-0.4073 - 0.0867 i)(J_5(r) + i Y_5(r) cos 5\emptyset$ + $(0.4077 - 0.8057 i)(J₆(r) + i Y₆(r) cos 6\emptyset)$ + $(0.5178 + 0.1445 i)(J₇(r) + i Y₇(r) cos 7 $\emptyset$$ + $(-0.0088 + 0.1321 i)(J_8(r) + i Y_8(r) \cos \theta)$ = $-0.7143 J_o(r) + 0.9634 J_1(r) \cos \beta + 0.7538 J_2(r) \cos 2\beta$ + 1.9084 $J_{z}(r)$ cos 3 \emptyset - 0.3627 $J_{\mu}(r)$ cos 4 \emptyset - 0.4073 $J_5(r)$ cos 5 \emptyset + 0.4077 $J_6(r)$ cos 6 \emptyset + 0.5178 $J_7(r)$ cos 7 \emptyset - 0.0088 $J_8(r)$ cos 8 \emptyset - 0.4517 $Y_0(r)$ + 0.7318 $Y_1(r)$ cos \emptyset + 0.9692 $Y_2(r)$ cos 2 \emptyset - 0.4178 $Y_7(r)$ cos 3 \emptyset + 0.7708 $Y_1(r)$ cos 4 \emptyset + 0.0867 $Y_5(r)$ cos 5 \emptyset + 0.8057 $Y_6(r)$ cos 6 \emptyset - 0.1445 $Y_7(r)$ cos 7 \emptyset - 0.1321 $Y_8(r)$ cos 8 \emptyset + i $\left[0.4517 \text{ J}_{0}(r) - 0.7318 \text{ J}_{1}(r) \cos\phi - 0.9692 \text{ J}_{2}(r) \cos 2\phi\right]$ + 0.4178 $J_3(r)$ cos 3 \emptyset - 0.7708 $J_4(r)$ cos 4 \emptyset - 0.0867 $J_5(r)$ cos 5 \emptyset - 0.8057 $J_6(r)$ cos 6 \emptyset + 0.1445 $J_7(r)$ cos 7 \emptyset + 0.1321 $J_8(r)$ cos 8 \emptyset - 0.7143 $Y_0(r)$ + 0.9634 $Y_1(r)$ cos \emptyset + 0.7538 $Y_2(r)$ cos 2 \emptyset + 1.9084 $Y_3(r)$ cos 3 \emptyset - 0.3627 $Y_{\mu}(r)$ cos 4 \emptyset - 0.4073 $Y_{5}(r)$ cos 5 \emptyset

Sists:

- + 0.4077 $Y_6(r)$ cos 6 \emptyset + 0.5178 $Y_7(r)$ cos 7 \emptyset
-

Consider the scattering of water waves from a vertical cylinder dipping through the water surface. Suppose a straight wave with wavelength equal to 2π is incident on the cylinder whose diameter is twelve units, which is almost two times the wavelength of incident wave. We want to see the outgoing and reflected waves near the cylinder. We shall approximate the scattered wave u_s in equation 10) by using the partial sum to m terms given by equation 14).

Let the partial sum for the scattered wave contain nine terms. Suppose that the wave is represented by the real part of equation 15), that is

Each term (T_m) and the partial sum (S_m) for $m = 1, 2, \ldots, 11$ are shown in figures 3.1 to 3.3 $\mathbf{T}_{\mathbf{m}}$ m $\mathtt{s}_{_{\mathrm{m}}}$ m Shows the graph of m^{th} term (T_m) and Figure 3.1 partial sum of m^{th} term (s_m) of R1(\mathcal{H}_s)_m for $m = 1, ..., 11, x = 0, y = 7.$ T_{m} ϕ $\mathtt{s}_{_{\mathrm{m}}}$ \mathbb{P} Shows the graph of m^{th} term (T_m) and Figure 3.2. partial sum of m^{th} term (S_m) of R1(\mathcal{H}_S)_m for $m = 1, ..., 11, x = 6 y = 7$ \mathbf{T}_m \mathbf{m} $\mathtt{s}_{_{\mathrm{m}}}$ m Figure 3.3. Shows the graph of m^{th} term (T_m) and partial sum of m^{th} term (S_m) of R1(\mathcal{H}_{S})_m for $m = 1, ..., 11, x = 8$ y = 7.

For the wave at time equal to zero the approximation to the complete wave along the x axis is

$$
\mathcal{H} = \text{RL} \left[e^{ikx} + (\mathcal{H}_s)_{8} \right]
$$

= $\cos x + R1$ (\mathcal{U}_s)₈, where k = 1.

Table 3.1 gives the values of R1 (\mathcal{U}_{s})₈ to four detimal places for $x = 0,1,2,..., 9$ and $y = 0,1,2,...,9$. Table 3.2 gives the values of R1 $(\mathcal{H}_s)_{8}$ to four decimal places for $x = 0, -1, -2, \ldots$, -9 and $y = 0,1,2,...,9$. Table 3.3 and 3.4 give the values of R1(\mathcal{H}_s) g+ cos x to four decimal places for $y = 0, 1, \ldots, 9, x = 0, 1, 2, \ldots$, and $x = 0, -1, \ldots, -9$, respectively

The diagram for cos x + R1 (γ_{s})₈ is shown in figures 3.4 and 3.5. Figure 3.4 shows the complete wave near a circular object. The incident wave is little affected by the scattering along the lines $y = \frac{1}{2}$ and y varies from left to right with the wavelength of the incident wave. On the left side of the object facing the incident wave the incident wave is combined with a large reflected wave, while behide the object the incident and scattered waves interfere destructively reducing the amplitude and creating a "Shadow". Figure 3.5 shows the contour curves of $\cos x + R1(\mathcal{H}_s)_{8}$.

ξ

 $\frac{3}{2}$

 $\frac{\nabla \tilde{\mathbf{F}}}{\tilde{\mathbf{Q}}\chi}$ 36

 $\mathbf x$ -7 -8 -9 -4 -5 -6 -3 -1 -2 \circ $\mathbf y$ $0.0744 - 0.8840$ -0.1200 1.1957 1.1313 \circ $-0.0054 - 0.9336$ 0.2761 1.2566 1.0484 $\mathbf 1$ $-0.2639 - 1.0329$ 0.8979 1.3147 0.8489 \overline{c} 0.8060 $-0.4845 - 1.0939$ 0.4750 1.0295 1.0568 $\overline{3}$ 0.4683 $-0.9682 - 1.0150$ -0.3228 0.4656 0.0506 0.0371 4 0.0684 -0.0353 -0.0838 $-0.5003 - 0.7478$ -0.9632 -1.0544 1.3482 -0.1945 1.1419 5 -0.0686 -0.3838 -0.9610 -0.0275 0.2131 -0.2585 -1.2772 -1.4527 1.0133 1.3273 6 0.8468 $0.4060 - 0.1383$ -0.4751 0.4930 -1.3686 -1.3717 1.4661 0.8824 -0.3704 $\overline{7}$ $0.7149 - 0.2294$ 0.1755 1.1369 1.4339 1.2541 0.7241 -0.4192 $-1.2007 - 0.9194$ $\rm 8$ $0.5261 - 0.6655$ -0.9655 -0.4956 1.6820 1.5885 0.7256 0.9856 -0.4368 0.5263 9

Table 3.4 Values of $(\frac{2}{5})_8$ + $\cos x$ to four decimal places.

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APPENDIX

EXPLANATION OF THE TABLES

The values in tables 2.1 to 2.5 in chapter II and tables 3.1 to 3.4 in chapter III are obtained by using tables 1 to 32.

Table 1 gives the values of r to two decimal places, where

$$
r = \sqrt{x^2 + y^2}, \quad x = 0,1,2, \ldots, 9
$$

$$
y = 0,1,2, \ldots, 9
$$

Tables 2to 10 give the values of the Bessel functions $J_n(r)$, $n = 0, 1, 2, ..., 8$ to four decimal places. The values of $J_0(r)$ and $J_1(r)$ are taken from "A Treatise on Bessel Functions and Their Applications to Physics" by A. Gray and T.M. Macrobert.

The values of $J_n(r)$ for n greater than one are obtained by using the recursion formula

 $J_{n+1}(r) = \frac{2n}{r} J_n(r) - J_{n-1}(r)$.

For example, find $J_2(1.41)$

Here $n = 1$, $r = 1.41$, $J_0(1.41) = 0.56142672$ and $J_1(1.41) = 0.54372550$.

Therefore $J_2(1.41) = \frac{2 \times 1}{1 \cdot 41} (0.54372550) - 0.56142672$

 $= 0.20981512$.

Tables 11 to 19 give the values of the Neumann functions $Y_n(r)$ of the first kind. Table 11 and table 12 give the values of $Y_0(r)$ and $Y_1(r)$ respectively calculated by using Bessel's interpolation formula.

These values must be calculated because the tables for $Y_0(r)$ and $Y_1(r)$ from "Theory of Bessel Functions" by Watson give values only for arguments having even numbers in the second decimal place, Tables'13 to 19 are obtained from the recursion formula above by substituting $Y_n(r)$ for $J_n(r)$.

Bessel's Interpolation

Suppose we are given values of a function $f(x)$ at equal intervals of x, say at x_0 , x_1 ,..., x_n where $x_r = x_0 + rh_r$ $r = 0$, ± 1 , ± 2 , ..., $\pm n$. Figure 1 shows the difference table for these values.

The interpolation at a point \bar{x} between x_0 and x_1 is obtained by using Bessel's formula

$$
f_{p} = f_{o} + p \left(f_{\frac{1}{2}} + B_{2} \left(\frac{d^{2}f_{o} + d^{2}f_{1}}{d^{2}} \right) + B_{3} \frac{d^{3}f_{1}}{d^{2}} \cdots \right)
$$
\nwhere $B_{2r} = \frac{1}{2} \left(\frac{p + r - 1}{2r} \right)$, $B_{2r+1} = \frac{p - \frac{1}{2}}{2r + 1} \left(\frac{p + r - 1}{2r} \right)$
\nand $p = \frac{\bar{x} - x_{o}}{h}$,

The array of differences involved is symmetric about a horizontal line midway between x_0 and x_1 .

The second differences are negligible if not greater than 4 , the third differences are negligible if less than 60, and the fourth differences are negligible if less than 20, where these values refer to the smallest significant units in the calculation.

For example, find $Y_{0}(2.23)$.

All values of $f(x)$ and their differences are multiplied by 10^{-7} .

Here
$$
p = \frac{2.23 - 2.22}{0.02}
$$
 = 0.5
\n
$$
B_2 = \frac{1}{4} p(p - 1)
$$
\n
$$
= \frac{1}{4} \times 0.5 (0.5 - 1)
$$
\n= -0.062.

The third difference is negligible since it is less than 60. Therefore

$$
Y_0(2.23) = 0.5206508 + 0.5 (-0.0003386)
$$

+ (-0.062)(-0.0002051 - 0.0002051)
= 0.52050693.

Tables 20 to 31 give the values of cos n \emptyset , n = 1, 2, ..., 8.

$$
\cos n \not\in = \frac{1}{2} \left[(2 \cos \beta)^n - \frac{n}{1} (2 \cos \beta)^{n-2} + \frac{n}{2} {n-3 \choose 1} (2 \cos \beta)^{n-4} - \frac{n}{3} {n-4 \choose 2} (2 \cos \beta)^{n-6} + \dots \right]
$$

Table 32 gives the values of cos x, x = 0, \pm 1, \pm 2,..., \pm 9.

$\mathbf x$ \mathbf{y}	\circ	ı	\overline{c}	3	4	5	6	7	8	9
\circ	0.00	1.00	2.00	3.00	4.00	5.00	6.00	7.00	8.00	9.00
ı	1.00	1.41	2.23	3.16	4.12	5.09	6.08	7.07	8.06	9.05
\overline{c}	2.00	2.23	2.82	3.60	4.47	5.38	6.32	7.28	8.24	9.21
3	3.00	3.16	3.60	4.24	5.00	5.83	6.70	7.61	8.54	9.48
4	4.00	4.12	4.47	5.00	5.65	6.40	7.21	8.06	8.94	9.84
5	5.00	5.09	5.38	5.83	6.40	7.07	7.81	8.60	9.43	10229
6	6.00	6.08	6.32	6.70	7.21	7.81	8.48	9.21	10.00	10.81
7	7.00	7.07	7.28	7.61	8.06	8.60	9.21	9.89	10.63	11.41
8	8.00	8.06	8.24	8.54	8.94	9.43	10.00	1063	11.31	12.04
9	9.00	9.05	9.21	9.48	9.84	10.29	10.81	11.41	12.04	12.72

Values of $r = x^2 + y^2$ to two decimal places ${\tt Table\ 1}$

 μ_{\downarrow}

x	\circ	ı	\overline{c}	3	4	5	6	7	8	9
\circ	1.0000	0.7652	0.2240		$-0.2601 - 0.3971 - 0.1776$		0.1506	0.3001	0.1717	-0.0903
ı	0.7652	0.5614	0.0937		-0.3094 -0.3865 -0.1477		0.1721	0.2997	0.1573	-0.1024
2	0.2240	0.0937	-0.1932	-0.3918		-0.3274 -0.0481	0.2279	0.2898	0.1118	-0.1389
3		$-0.2601 - 0.3094$	-0.3918	-0.3707	-0.1776	0.1010	0.2851	0.2500	0.0310	-0.1907
4	-0.3971	-0.3865	-0.3274	-0.1776	0.0436	0.2433	0.2945		$0.1573 - 0.0754$	-0.2358
5°	-0.1776	-0.1477	-0.0481	0.1010	0.2433	0.2997	0.2134		0.0146 -0.1821	-0.2480
6	0.1506	0.1721	0.2279	0.2851	0.2945	0.2134	0.0474	-0.1389 -0.2459		-0.2018
7	0.3001	0.2997	0.2898	0.2500	0.1573		0.0146 -0.1389	-0.2396 -0.2245		-0.0880
8	0.1717	0.1573	0.1118	0.0310		-0.0754 -0.1821 -0.2459		-0.2245 -0.1090		0.0566
9	-0.0903 -0.1024		-0.1389	-0.1907	-0.2358 -0.2480 -0.2018			-0.0880	0.0566	0.1792

Table 2. Values of $J_0(r)$ to four decimal places

 $\mathbf{x} \approx \mathbf{z}^{\mathbf{R}^{\mathrm{U}}} \mathbf{y}^{\mathbf{q}^{\mathrm{U}}} = \mathbf{y}$

 $\sharp\tilde{5}$

Table 4. Values of $J_2(r)$ to four decimal places

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x \mathbf{y}	\circ	ı	2	$\overline{3}$	4	5	6	7	8	9
\circ	0.0000	0.0002	0.0070	0.0430	0.1321	0.2611	0.3621	0.3479	0.1858	-0.0550
ı	0.0002	0.0013	0.0116	0.0534	0.1463	0.2727	0.3662	0.3412	0.1723	-0.0667
\overline{c}	0.0070	0.0116	0.0331	0.0897	0.1908	0.3078	0.3734	0.3164	0.1305	-0.1027
3	0.0430	0.0534	0.0897	0.1611	0.2611	0.3509	0.3680	0.2645	0.0572	-0.1576
4	0.1321	0.1463	0.1908	0.2611	0.3358	0.3741	0.3254	0.1723	-0.0408	-0.2153
5	0.2611	0.2727	0.3078	0.3509	0.3741	0.3412	0.2262	0.0424	-0.1481	-0.2558
6	0.3621	0.3662	0.3734	0.3680	0.3254	0.2262	0.0721	-0.1027	-0.2341	-0.2527
7	0.3479	0.3412	0.3164	0.2645	0.1723	0.0424	-0.1027	-0.2217	-0.2599	-0.1860
8	0.1858	0.1723	0.1305	0.0572	-0.0408	-0.1481	-0.2341	-0.2599	-0.2011	-0.0649
9	-0.0550	-0.0667	-0.1027	-0.1576	-0.2153	-0.2558	-0.2527	-0.1860	-0.0649	0.0800

Table 7: Values of $J_5(r)$ to four decimal places

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$\mathbf x$ У	\circ	ı	\overline{c}	3	4	5	6	7	8	9
\circ	0.0000	0.0000	0.0012	0.0114	0.0491	0.1310	0.2458	0.3392	0.3376	0.2043
ı	0.0000	0.0002	0.0022	0.0150	0.0564	0.1405	0.2551	0.3430	0.3333	0.1946
\overline{c}	0.0012	0.0022	0.0082	0.0293	0.0818	0.1728	0.2816	0.3511	0.3174	0.1623
3	0.0114	0.0150	0.0293	0.0645	0.1310	0.2259	0.3180	0.3534	0.2812	0.1038
4	0.0491	0.0564	0.0818	0.1310	0.2045	0.2900	0.3489	0.3333	0.2156	0.0219
5	0.1310	0.1405	0.1728	0.2259	0.2900	0.3430	0.3479	0.2725	0.1149	-0.0779
6	0.2458	0.2551	0.2816	0.3180	0.3489	0.3479	0.2894	0.1623	-0.0145	-0.1745
7	0.3392	0.3430	0.3511	0.3534	0.3333	0.2725	0.1623	0.0105	-0.1444	-0.2405
8	0.3376	0.3333	0.3174	0.2812	0.2156	0.1149	-0.0145	-0.1444	-0.2336	-0.2416
9	0.2043	0.1946	0.1623	0.1038	0.0219	-0.0779	-0.1745	-0.2405	-0.2416	-0.1665

Table 8: Values of $J_6(r)$ to four decimal places

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Table 12: Values of $Y_n(r)$ to four decimal places

Table 13: Values of $Y_2(r)$ to four decimal places

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Table 21 : Value of $cos \phi$ to four decimal places.

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Table 22 : Value of $cos 2 \notimes 1$ to four decimal places

Table 23 : Values of $cos 3 \not\theta$ to four decimal places.

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Table 24 : Values of cos $\frac{1}{2}$ \emptyset to four decimal places.

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x y	\circ	-1	-2,	-3	-4	-5	-6	-7	-8	-9
\circ	0.0000	-1.0000	-1.0000	-1.0000	-1.0000	-1.0000	-1.0000	-1.0000	-1.0000	-1.0000
ı	0.0000	0.7175	0.6600	0.0282	-0.3518	-0.5881	-0.6871	-0.7599	-0.8181	$ -0.8631$
\overline{c}	0.0000	-0.7289	0.7175	0.9774	0.6767	0.3152	0.0282	-0.1779	-0.3518	-0.4783
3	0.0000	-0.9992	-0.1955	0.7091	0.9971	0.9045	0.6709	0.4308	0.2161	0.0282
4	0.0000	-0.9411	-0.7327	0.0758	0.7111	0.9738	0.9797	0.8539	0.6767	0.4887
5	0.0000	-0.8352	-0.9449	-0.4255	0.2319	0.7076	0.9453	0.9991	0.9389	0.8193
6	0.0000	$-C.7351$	-0.9992	-0.7316	-0.2002	0.3261	0.7091	0.9227	0.9971	0.9789
$\overline{7}$	0.0000	-0.6514	-0.9839	-0.8982	-0.5139	-0.0400	0.3945	0.7101	0.8991	0.9858
$\,$ 8	0.0000	-0.5823	-0.9411	-0.9751	-0.7327	-0.3405	0.0758	0.4372	0.7081	0.8819
9	0.0000	-0.5254	-0.8886	-0.9992	-0.8667	-0.5685	-0.1985	0.1595	0.4723	0.7091

Table 27 : Values of cos $5 \not\! 0$ to four decimal places

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