## CHAPTER III

### FUNDAMENTAL OF DISTILLATION

### 3.1 Theoretical Model for Multicomponent Multistage Distillation Column.

The theoretical model for this thesis has the assumptions as follows :

- This column is continuous distillation.
- It is operated at steady-state conditions.
- The liquid and vapor flows are countercurrent.
- No chemical reaction occurs.
- Phase equilibrium is achieved at each stage.
- The stages are numbered down from condenser to rebeiler.
- Each stage has feed, vapor-liquid interstage streams, vapor-liquid sidestreams and heat transfer that are shown in figure 3.1.

Considering stage j in figure 3.1. The single or two phase feed enter stage j with flow rate  $F_j$  which contains overall composition in mole fractions  $z_{i,j}$  of component i, temperature  $T_{Fj'}$  pressure  $P_{Fj}$  and overall enthalpy  $H_{Fj}$ .

The interstage liquid stream from stage above enters stage j with flow rate  $L_{j-1}$  contains the liquid composition in mole fractions  $x_{i,j-1}$  of component i, temperature  $T_{j-1}$ , enthalpy  $H_{Lj-1}$  and pressure  $P_{j-1}$  which is equal to or less than the pressure of stage j.

Similarly, the interstage vapor liquid stream from stage j+1 below flows into stage j which has flow rate  $V_{i+1}$ , the vapor composition in mole fraction  $y_{i+1}$ , temperature  $T_{j+1},$  enthalpy  $H_{Vj+1}$  and pressure  $P_{j+1}$  that is equal to or higher than  $P_{j}.$ 



Vapor from stage below

Figure 3.1 Equilibrium Stage

The vapor and liquid leaving stage j have the intensive properties  $P_j$  and  $T_j$ . The vapor can be divided into vapor sidestream of flow rate  $W_j$  and the interstage vapor of flow rate  $V_j$ , with the composition of component i  $y_{i,j}$  and enthalpy  $H_{V_j}$ . The liquid stream is similar to vapor stream that has liquid

sidestream  $U_{j},$  interstage liquid flow rate  $L_{j},$  liquid composition  $x_{i,j}$  and enthalpy  $H_{Lj}.$ 

In this system, the C components mixture is separated by N stages distillation column. The distillation calculation can be described by the set of equations that are defined from four relations. (Henry, E. J., 1981)

1. Material balance equation for component i on stage j.

$$L_{j-1} x_{i,j-1} + V_{j+1} y_{i,j+1} + F_j z_{i,j} - (L_j + U_j) x_{i,j} - (V_j + W_j) y_{i,j} = 0$$
(3-1)

2. Phase equilibrium relation component i on stage j.

$$y_{i,j} - K_{i,j} x_{i,j} = 0$$
 (3-2)

where  $\boldsymbol{K}_{i,j}$  is the phase equilibrium constant of component i at stage j.

3. Mole fraction summation on stage j.

$$\sum_{i=1}^{C} y_{i,i} - 1.0 = 0$$
(3-3)

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$$\sum_{i=1}^{C} \mathbf{x}_{i,i} - 1.0 = 0 \tag{3-4}$$

4. Energy balance on stage j.

$$L_{j-1}H_{L_{j-1}} + V_{j+1}H_{V_{j+1}} + F_{j}H_{F_{j}} - (L_{j}+U_{j})H_{L_{j}} - (V_{j}+W_{j})H_{V_{j}} - Q_{j} = 0$$
(3-5)

where kinetic and potential energy changes are assumed to be insignificant.

 $\rm H_{L},~\rm H_{V},~and~\rm H_{F}$  are mixture enthalpy of the corresponding stream.

A total material balance equation can be used in place of equation (3-3) or (3-4). It is derived by combining these two equations and  $\sum_{i} z_{i} = 1.0$  with (3-1) summed over the C components and over stages 1 through j to yield

$$L_j = V_{j+1} + \sum_{m=1}^{j} (F_m - W_m - U_m) - V_1$$
  $j = 1, 2, 3, ..., N$  (3-6)

In general, it is known that  $K_{i,j} = K_{i,j} \{T_j, P_j, x_{i,j}, y_{j,j}\}, H_{Vj} = H_{Vj} \{T_j, P_j, y_{j,j}\}$ , and  $H_{Lj} = H_{Lj} \{T_j, P_j, x_{j,j}\}$ . If these relations are not counted as equations and three properties are not counted as variables, this column has 2C+3 equations for each equilibrium stage. If N and all  $F_j$ ,  $z_{i,j}$ ,  $T_{Fj}$ ,  $P_{Fj}$ ,  $P_j$ ,  $U_j$ ,  $W_j$  and  $O_j$  are specified, the model is represented by N(2C+3) simultaneous algebraic equations in N(2C+3) unknown. (Shuzo Ohe, 2536) This model has a large of nonlinear equations that must be solved by iterative techniques.

#### 3.2 Tridiagonal Matrix Algorithm.

There are many iterative solution precedures for solving nonlinear algebraic equations. For separators where the feeds centains only components of similar volatility (narrow-boiling case), the bubble-point (BP) method is recommended. The key to the success of BP method is the tridiagonal matrix that is resulted from a modified material balance equations (3-1) when  $T_i$  and  $V_j$  are selected as the tear variables. The material balance equations are modified in the unknown liquid mole fractions. This set of equations for each component is

solved by highly efficient and reliable algorithm due to Thomas as applied by Wang and Henke.

The  $y_{i,j}$  in equation (3-1) are substituted by equation (3-2). The material balance equations are obtained by substituting equation (3-6) into equation (3-1) to eliminate  $L_j$ . The result for each component and each stage is shown as equation (3-7).

$$A_{j}x_{i,j+1} + B_{i,j}x_{i,j} + C_{i,j}x_{i,j+1} = D_{i,j}$$
 (3-7)

-

where

$$A_j = V_j + \sum_{m=1}^{j-1} (F_m - W_m - U_m) - V_1 \qquad j = 2, 3, 4, ..., N$$
 (3-8)

$$B_{i,j} = - [V_{j+1} + \sum_{m=1}^{j} (F_m - W_m - U_m) - V_1 + U_j + (V_j + W_j) K_{i,j}]$$

$$j = 1, 2, ..., N$$

$$i = 1, 2, ..., C$$
(3-9)

$$C_{i,j} = V_{j+1} K_{i,j+1}$$
  $j = 1, 2, 3, ..., N-1$  (3-10)  
 $i = 1, 2, 3, ..., C$ 

$$D_{i,j} = -F_j z_{i,j} \qquad j = 1, 2, 3, ..., N \qquad (3-11)$$
  
$$i = 1, 2, 3, ..., C$$

with  $\boldsymbol{x}_{i,0}$  = 0,  $\boldsymbol{V}_{N+1}$  = 0,  $\boldsymbol{W}_1$  = 0, and  $\boldsymbol{U}_N$  = 0.

If the modified equations are grouped by component, they can be represented by tridiagonal matrix equations (3-12).

$$\begin{bmatrix} B_{1} & C_{1} & 0 & 0 & \cdots & 0 & 0 \\ A_{2} & B_{2} & C_{2} & 0 & \cdots & 0 & 0 \\ 0 & A_{3} & B_{3} & C_{3} & \cdots & 0 & 0 \\ 0 & 0 & \cdots & \cdots & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & \cdots & \cdots & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & \cdots & \cdots & 0 & A_{N-1} & B_{N-1} & C_{N} \\ 0 & 0 & \cdots & \cdots & 0 & A_{N-1} & B_{N} \end{bmatrix} \begin{bmatrix} x_{1,1} \\ x_{1,2} \\ \vdots \\ x_{1,3} \\ \vdots \\ \vdots \\ x_{1,N-1} \\ x_{1,N} \end{bmatrix} = \begin{bmatrix} D_{1} \\ D_{2} \\ D_{3} \\ \vdots \\ D_{3} \\ \vdots \\ D_{N-1} \\ D_{N} \end{bmatrix}$$
(3-12)

The Thomas algorithm for solving the equation (3-12) is a Gaussian eliminating  $\mathbf{x}_{i,j}$ . Firstly, the forward elimination starts from stage 1 and works toward stage N to find  $\mathbf{x}_{i,N}$ . Then, the other  $\mathbf{x}_{i,j}$  are obtained starting with  $\mathbf{x}_{i,N-1}$  by backward substitution.

The equations used in the Thomas algorithm are as follows:

$$p_1 = C_1 / B_1$$
 (3-13)

$$q_1 = D_1 / B_1$$
 (3-14)

$$p_j = \frac{C_j}{(B_j - A_j p_{j-1})}$$
  $j = 2, 3, 4, ..., N-1$  (3-15)

$$q_j = \frac{(D_j - A_j q_{j-1})}{(B_j - A_j p_{j-1})}$$
 j = 2, 3, 4, ..., N (3-16)

The equation (3-7) is rearranged as

$$\mathbf{x}_{i,j} = \mathbf{q}_j - \mathbf{p}_j \mathbf{x}_{i,j+1}$$
 (3-17)

The forward elimination starts from stage 1 to stage N. For stage N, (3-17) isolate  $x_{i,N}$  as

$$\mathbf{x}_{1,\mathbf{N}} = \mathbf{q}_{\mathbf{N}} \tag{3-18}$$

Then, the backward substitution eliminates the other  $x_{i,j}$  by equation (3-19).

$$\mathbf{x}_{i,j} = \mathbf{q}_j - \mathbf{p}_j \mathbf{x}_{i,j+1}$$
 j = N-1, N-2, ..., 1 (3-19)

In general, computed values of  $\boldsymbol{x}_{i,j}$  are always positive.

## 3.3 The Bubble-Point (BP) Method for Distillation.

The effective solution procedure for this work is referred to the bubblepoint (BP) method because a new set of stage temperatures is computed during each iteration from bubble-point equation.

Computational procedure:

Step 1. Specify feed and column configuration as follows:

- All of feeds and feed conditions ( $F_{j}$ ,  $z_{ij}$ ,  $T_{Fj}$ ,  $P_{Fj}$  or  $H_{Fj}$ ).
- Column configurations N, L (reflux rate),  $V_1$  (vapor distillate rate),  $P_j$ ,  $U_j$ ,  $W_j$ , and  $Q_j$  except  $Q_1$  (condenser duty) and  $Q_N$  (reboiler duty).

Step 2. Initialize tear variables  $\boldsymbol{T}_{j}$  and  $\boldsymbol{V}_{j}.$ 

- To initiate a set of  $V_j$  values based on the assumption of constant interstage flows.
- To initiate a set of Tj values by assuming a linear variation of temperatures with stage location.

Step 3. Solve equation (3-12) for each set of  $x_{i,j}$  values by the Thomas algorithm.

Step 4. Normalize the set of  $x_{i,j}$  values for each stage by the relation

$$(\mathbf{x}_{i,j})_{\text{normalized}} = \frac{\mathbf{x}_{i,j}}{\sum_{i=1}^{c} \mathbf{x}_{i,j}}$$
(3-20)

Step 5. Compute the new set of  $T_{\rm j}$  from bubble point equation (3-21) and set of  $y_{\rm i,j}$  values from (3-2).

$$\sum_{i=1}^{c} K_{ij} x_{ij} - 1.0 = 0$$
(3-21)

The equation (3-21) is nonlinear. It must be solved by iterative method.

Step 6. Compute condenser duty from (3-5) and reboiler duty from (3-22).

$$Q_{N} = \sum_{j=1}^{N} (F_{j}H_{Fj} - U_{j}H_{Lj} - W_{j}H_{Vj}) - \sum_{j=1}^{N-1} Q_{j} - V_{1}H_{V1} - L_{N}H_{LN}$$
(3-22)

Step 7. Compute the new set of  $V_j$  from modified energy balance equation, which is obtained by combining equations (3-5) and (3-6) twice to eliminate  $L_{j-1}$  and  $L_j$ .

$$\alpha_{j}V_{j} + \beta_{j}V_{j+1} = \gamma_{j}$$
(3-23)

where

$$\alpha_{j} = H_{Lj-1} - H_{vj}$$
(3-24)

$$\beta_{j} = H_{V_{j+1}} - H_{l_{j}}$$
 (3-25)

$$\gamma_{j} = \left[\sum_{m=1}^{j-1} (F_{m} - W_{m} - U_{m}) - V_{1}\right](H_{L_{j}} - H_{L_{j-1}}) + F_{j}(H_{L_{j}} - H_{F_{j}}) + W_{j}(H_{V_{j}} - H_{L_{j}}) + O_{j}$$
(3-26)

The modified energy balance is written in didiagonal matrix from. It is applied over stages 2 to N-1 as follow:

$\beta_2$	0	0	0		-	0	0	V 3		$\gamma_2 - \alpha_2 V_2$	
a	etas	0	0	÷	÷	0	0	V 4		<i>γ</i> 3	
0	04	$\beta_{4}$	0		-	0	0	V 5		Y4	
0	0	-				0	0	-	ļ		
-		÷		-	-	-	-	-	=		(3-27)
-	÷	÷	-	÷	-		-	1.1		(1.5)	
0	0	-	-		<i>Q</i> N 2	$\beta_{\rm N}$ 2	0	V N-I		γN-2	
0	0	-	-		0	<b>a</b> n 1	βN 1	V N		γN-1	

Using Gaussian elimination procedure to solve equation (3-27) from the top where  $V_2$  is computed by equation (3-6).

$$V_{j} = \frac{\gamma_{j-1} - \alpha_{j-1} V_{j-1}}{\beta_{j-1}}$$
(3-28)

Then, compute the set of  $L_1$  from equation (3-6). Step 8. Check the convergence of this procedure.

$$\sum_{j=1}^{N} \left[ \frac{T_{j}^{(k)} - T_{j}^{(k-1)}}{T_{j}} \right]^{2} + \sum_{j=1}^{N} \left[ \frac{V_{j}^{(k)} - V_{j}^{(k-1)}}{V_{j}} \right]^{2} \leq \varepsilon$$
(3-29)

where k is the iteration index.

 $\boldsymbol{\epsilon}$  is prescribed tolerance.

However, Wong and Henke suggested that the following simpler criterion which is based on successive sets of Tj values only is adequate.

$$\tau = \sum_{j=1}^{N} \left[ T_{j}^{(k)} - T_{j}^{(k-1)} \right]^{2} \le 0.01 \text{N}$$
(3-30)

and

$$\tau = \sum_{j=1}^{N} \left[ V_{j}^{(k)} - V_{j}^{(k-1)} \right]^{2} \le 0.01 \text{N}$$
(3-31)

The new set of  $\boldsymbol{V}_{j}$  and  $\boldsymbol{T}_{j}$  values change 10% from last iterations.

Step 9. Use the new set of Vj and Tj as the initial guesses for calculating step 3 to 9 when the result of step 8 is false. Stop the calculation when which result become true.

The algorithm for the Wang-Henke BP method is shown in figure 3.2.



Figure 3.2 Algorithm for Wang-Henke BP method (Henry, E. J., 1981)

## 3.4 Vapor Liquid Equilibria (VLE)

The liquid mixture at temperature T and pressure P is in equilibria with vapor mixture at the same temperature and pressure. (Van Ness, H. C., 1982) The condition of thermodynamic equilibria for every component i in the mixture is given by

$$\bar{f_i^{\nu}} = \bar{f_i^{\,L}} \tag{3-31}$$

The Raoult's Law shown in equation (3-32) is applied for VLE of the ideal mixture that the liquid phase is an ideal solution, vapor phase is an ideal gas, and the liquid phase fugacities are independent of pressure.

$$y_i P = x_i P_i$$
(3-32)

Vapor liquid equilibria are often expressed in term of phase equilibria ratio.

$$K_{i,j} = \frac{Y_{i,j}}{X_{i,j}}$$
(3-33)

K values for ideal mixture are defined as

$$K_{i,j} = \frac{P_{i,j}^{\bullet}}{P}$$
(3-34)

For vapor liquid equilibria at high pressure are conveniently calculated by using an equation of state applicable to both phases. Equation (3-33) can be rewritten in terms of coefficients which defined by

$$\phi_i^{\mathsf{v}} = \frac{\mathbf{f}_i^{\mathsf{v}}}{\mathbf{y}_i \mathbf{P}} \tag{3-35}$$

and

$$\phi_i^L = \frac{\mathbf{f}_i^L}{\mathbf{x}_i P} \tag{3-36}$$

Phase equilibrium ratio is derived by combining equation (3-35) and (3-36) into equation (3-31) to give

$$K_i = \frac{\phi_i^{\prime}}{\phi_i^{\prime}} \tag{3-37}$$

Assuming that the equation of state, can be expressed as

$$P = f(T, V, z_1, z_2, ...)$$
(3-38)

And assuming further that this equation of state holds for all fluid densities (i. e., gases and liquids) and for all compositions z1, z2, ..., we can calculate  $\phi_i^L$  and  $\phi_i^V$  from

$$\operatorname{RT} \ln \phi_i^L = \int_{V^L}^{\infty} \left[ \left( \frac{\partial P}{\partial n_i} \right)_{T,V,n_j} - \frac{RT}{V} \right] dV - \operatorname{RT} \ln Z^L$$
(3-39)

$$\operatorname{RT} \ln \phi_{i}^{\mathsf{v}} = \int_{\mathsf{v}^{\mathsf{v}}}^{\infty} \left[ \left( \frac{\partial \mathsf{P}}{\partial \mathsf{n}_{i}} \right)_{\mathsf{T},\mathsf{v},\mathsf{n}_{i}} - \frac{\mathrm{RT}}{\mathrm{V}} \right] \, \mathrm{dV} - \operatorname{RT} \ln \mathsf{Z}^{\mathsf{V}}$$
(3-40)

where compressibility factor Z is given by

$$Z^{L} = \frac{PV^{L}}{RT}$$
(3-41)

$$Z^{V} = \frac{PV^{V}}{RT}$$
(3-42)

. .

In the liquid phase, the total volume  $V_{T}^{\text{L}}$  is related to the molar volume  $V_{}^{\text{L}}$ 

$$V^{L} = \frac{V_{T}^{L}}{n_{T}^{L}}$$
(3-43)

where  $n_T^L$  is the total number of moles in the liquid phase. Similarly,

$$V^{V} = \frac{V_{T}^{V}}{n_{T}^{V}}$$
(3-44)

In principle. Equations (3-37) to (3-38) are sufficient for finding all K factors in a multicomponent system containing two fluid phases. However, if a realistic equation of state is used, the required computations are strongly nonlinear and often require extensive iterations.

#### 3.5 Bubble Point Calculation

by

The bubble point calculation used for computing a new set of temperatures is very important step of BP method. It is particularly effective for mixtures having a narrow range of K values because temperature is not then sensitive to composition.

The bubble point criterion is

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$$1 = \sum_{i=1}^{c} z_{i} K_{i}$$
 (3-45)

with  $x_i = z_i$  and  $y_i = K_i x_i$ 

Equation (3-45) is used for calculating bubble point temperature at a given pressure. It is highly nonlinear in temperature. Therefore, iterative procedures are required to solve bubble point temperatures. Calculation procedure for ideal case is shown in figure 3.3. (Henry, E. J., 1981)



Figure 3.3 An algorithm for solving the bubble point temperature.

The bubble point calculation using equations of state is shown in figure 3.4.



Figure 3.4 An algorithm for solving the bubble point temperature using equation of state. (Sandler, S. I., 1989)

# 3.6 The thermodynamices Properties

The thermodynamic properties presented in this chapter are widely used for distillation calculation

3.6.1 Vapor Pressure

The widely used vapor pressure correlation is the Antoine equation

$$\ln P_i^* = A_i - \frac{B_i}{T + C_i}$$
(3-46)

1

where  $P_i^* = Vapor Pressure$ 

 $\begin{array}{ll} T & = Temperature \\ A_{i^{\prime}} & B_{i} \mbox{ and } C_{i} = \mbox{constant parameters} \end{array}$ 

The applicable temperature range is not large and corresponds to a pressure interval of about 0.01 to 2 bar. The Antoine equation should not be used outside the temperature limits stated. (Reid, R. C., 1988)

The other correlation predicting vapor pressure over wide range of temperature is the Wagner's equation.

$$\ln \left[\frac{P^{*}}{P_{c}}\right] = \frac{a \tau + b \tau^{1.5} + c \tau^{3} + d \tau^{6}}{Tr}$$
(3-47)

where  $\tau = 1 - Tr$  $Tr = \frac{T}{Tc}$  The enthalpy used in energy balance equation is very important for solving distillation problems.

The molal specific heat of gases is conventionally given as a polynomial in temperatures. The superscript <sup>o</sup> refers to the ideal condition.

$$\mathbf{c}_{Pv}^{o} = \mathbf{a}_{1} + \mathbf{a}_{2}\mathbf{T} + \mathbf{a}_{3}\mathbf{T}^{2} + \mathbf{a}_{4}\mathbf{T}^{3} + \mathbf{a}_{5}\mathbf{T}^{4}$$
(3-48)

. .

The integral of equation (3-48) between a referent temperature,  $T_0$ , to a desired temperture, T, provides ideal gas molar enthalpy at that temperature, as below:

$$H_{v}^{o} = \int_{T_{o}}^{T} c_{Pv}^{o} dT = \sum_{k=1}^{5} \frac{a_{k} (T^{k} - T_{o}^{k})}{k}$$
(3-49)

For real gas the estimated expression of enthalpy must be modified to yeild

$$(H-H_{v}^{o}) = Pv - RT - \int_{\infty}^{v} \left[P - T\left(\frac{\partial P}{\partial T}\right)_{v}\right] dv$$
(3-50)

where V is the total volume equal to  $v \sum n_i$ .

is

For a mixture at given temperature T and pressure P, the vapor enthalpy

$$H_{v} = \left[\sum_{i=1}^{c} (y_{i} H_{iv}^{o})\right] + (H_{v} - H_{v}^{o})$$
(3-51)

Equations (3-48) and (3-50) particularly suitable by using equations of state are explicit in pressure.

$$H_{L} = \left[\sum_{i=1}^{c} (y_{i} H_{iv}^{o})\right] + (H_{L} - H_{v}^{o})$$
(3-52)

The equation of state mentianed above also does not apply with polar compound, large molecules (polymer) and electrolytes. (Thanit Sawasdisevi, 1996)